

## 5. Initial value problem - Homogeneous medium

Deep water waves ( $\omega^2 = gk$ ). Again method of stationary phase

$$\text{Put } \eta(x,t) = \int_{-\infty}^{+\infty} [C(k)e^{i(kx+\omega t)} + D(k)e^{i(kx-\omega t)}] dk$$

superposition of waves going in opposite directions

We need to specify at  $t = 0$

$$\eta(x,0) = \int_{-\infty}^{+\infty} [C(k) + D(k)] e^{ikx} dk$$

and

$$\eta_t(x,0) = \int_{-\infty}^{+\infty} i\omega [C(k) - D(k)] e^{ikx} dk$$

Assume for simplicity  $\eta_t(x,0) = 0$

then  $C(k) = D(k)$  and  $C(k) + D(k) = 2C(k) = \tilde{\eta}_0$

Then

$$\eta(x,t) = \frac{1}{2} \int_{-\infty}^{+\infty} \tilde{\eta}_0(k) \left[ e^{i(kx+\omega t)} + e^{i(kx-\omega t)} \right] dk$$

Now  $\omega = (\text{sign } k)(g|k|)^{1/2}$  to ensure that waves at frequency  $\omega$  propagate in the same direction (even though oppositely) regardless of the sign of  $k$ . Separate into left and right-going wave contributions:

$$\eta(x,t) = \frac{1}{2} \int_{-\infty}^{+\infty} \tilde{\eta}_0(k) e^{i\theta^+ t} dk + \frac{1}{2} \int_{-\infty}^{+\infty} \tilde{\eta}_0(k) e^{i\theta^- t} dk$$

where

$$\theta^\pm \equiv \frac{kx}{t} \pm (\text{sign } k)(g|k|)^{1/2}$$

Points of stationary phase (giving the major contribution to the integrals) are those where

$$\frac{\partial}{\partial k} \theta^\pm = 0$$

Let us consider the half-plane ( $x > 0$ ) for  $k > 0$   $|k| = R$

$$\theta^+(k) = \frac{kx}{t} + (gk)^{1/2} \quad \text{left-going wave}$$

$$\frac{\partial \theta^+}{\partial k} = \frac{x}{t} + \frac{1}{2} \sqrt{\frac{g}{k}} = 0$$

$$\text{gives } x = -\frac{t}{2} \sqrt{\frac{g}{k}} < 0 \quad \text{always for increasing time}$$

No stationary points in  $\theta^+$  for  $x > 0$

$$\theta^- = \frac{kx}{t} - \frac{1}{2}(gk)^{1/2} = 0 \quad \text{right going wave}$$

$$\frac{\partial \theta^-}{\partial k} \Big|_{k=k_0} = \frac{x}{t} - \frac{1}{2} \left(\frac{g}{k}\right)^{1/2} = 0 \quad \text{at } k = k_0$$

$$\frac{x}{t} - \frac{1}{2} \left(\frac{g}{k_0}\right)^{1/2} = c_g \Big|_{k_0} \quad \text{group velocity of packet centered at } k_0$$

$$\frac{\partial^2 \theta^-}{\partial k^2} \Big|_{k_0} = \frac{1}{4} \left(\frac{g^{1/2}}{k_0^{3/2}}\right) = \frac{2x^3}{gt^3} \quad \text{as } \frac{1}{k_0^{3/2}} = \frac{8x^3}{t^3} \frac{1}{g^{3/2}}$$

only right going wave

$$\text{Then } \eta_{k>0}(x > 0, t) \simeq \frac{1}{2} \tilde{\eta}_0(k_0) e^{i\theta^-(k_0)t} \int_{-\infty}^{+\infty} e^{\frac{(k-k_0)^2 \theta''(k_0)t}{2i}} dk$$

$$\text{As } \int_{-\infty}^{+\infty} e^{-\alpha z^2} dz = \left(\frac{\pi}{\alpha}\right)^{1/2} \quad \text{and } z^2 = (k-k_0)^2$$

$$\alpha = \frac{\theta''(k_0)t}{2i}$$

$$\eta_{k>0}(x > 0, t) \simeq \frac{1}{2} \tilde{\eta}_0(k_0) e^{i\theta^-(k_0)t} \left[ \frac{2\pi i}{t\theta''(k_0)} \right]^{1/2}$$

$$\text{But } \frac{2\pi i}{t\theta''(k_0)} = e^{i\frac{\pi}{2} \left( \frac{\pi g t^2}{x^3} \right)}$$

Hence

$$\theta^-(k_o) = \left[ k_o \frac{x}{t} - g^{1/2} k_o^{1/2} \right] t$$

$$= \left[ k_o \frac{g^{1/2}}{2k_o^{1/2}} - g^{1/2} k_o^{1/2} \right] t = -\frac{1}{2} g^{1/2} k_o^{1/2} t$$

As  $k_o^{3/2} = \frac{g^{3/2} t^3}{8x^3} \Rightarrow k_o^{1/2} = \frac{g^{1/2} t}{2x}$

And

$$\theta^-(t_o) = -\frac{1}{2} g^{1/2} t \frac{g^{1/2} t}{2x} = -\frac{1}{4} \frac{gt^2}{x}$$

So  $e^{i\theta^-(k_o)t} \left[ \frac{2\pi i}{t\theta''(k_o)} \right]^{1/2} = e^{-\frac{1}{4} i \frac{gt^2}{x}} e^{i\frac{\pi}{4} \left[ \frac{\pi gt^2}{x^3} \right]^{1/2}} = e^{-i \left[ \frac{gt^2}{4x} - \frac{\pi}{4} \right] \left[ \frac{\pi gt^2}{x^3} \right]^{1/2}}$

Take the real part

$$\eta_{k>o}(x > o, t) \cong \frac{1}{2} \tilde{\eta}_o(k_o) \left[ \frac{\pi gt^2}{x^3} \right]^{1/2} \cos \left[ \frac{gt^2}{4x} - \frac{\pi}{4} \right]$$

$$\sim \frac{1}{2} \tilde{\eta}_o(k_o) (\pi g)^{1/2} \left( \frac{t}{x^{3/2}} \right) \cos \left[ \frac{gt^2}{4x} - \frac{\pi}{4} \right]$$

$$\left[ \frac{\pi gt^2}{x^3} \right]^{1/2} \text{ is the modulating amplitude}$$

Central wavelength  $\frac{2\pi}{k_o} = \frac{8\pi x^2}{gt^2}$

$x \rightarrow$  increases  $\lambda \rightarrow$  increases at fixed  $x$

The final solution is:

$$\eta(x > 0, t) \approx \frac{1}{2} \hat{\eta}_0(k_0) (\pi g)^{1/2} \left( \frac{t}{x^{3/2}} \right) \cos \left[ \frac{gt^2}{4x} - \frac{\pi}{4} \right]$$

modulating wave  
amplitude part

Plot  $\eta(x, t)$  as a function of  $x$  at fixed  $t$ , i.e. a snapshot of the wave:

The central wavelength  $\frac{2\pi}{k_0}$  increases and the amplitude decreases with  $x$  going away from the initial position.

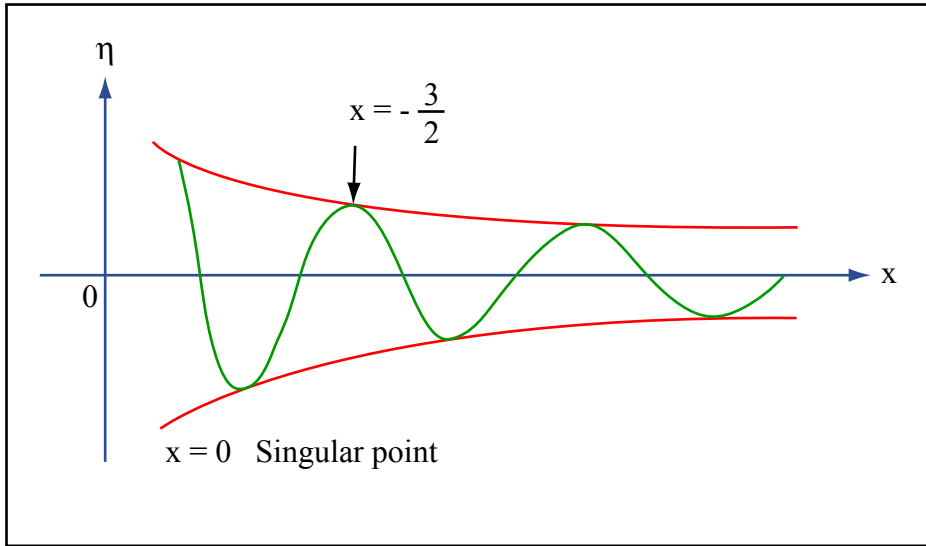
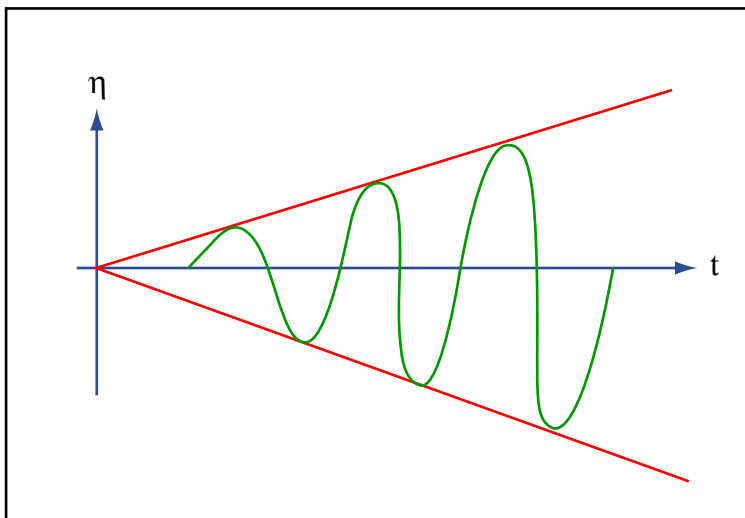


Figure by MIT OpenCourseWare.

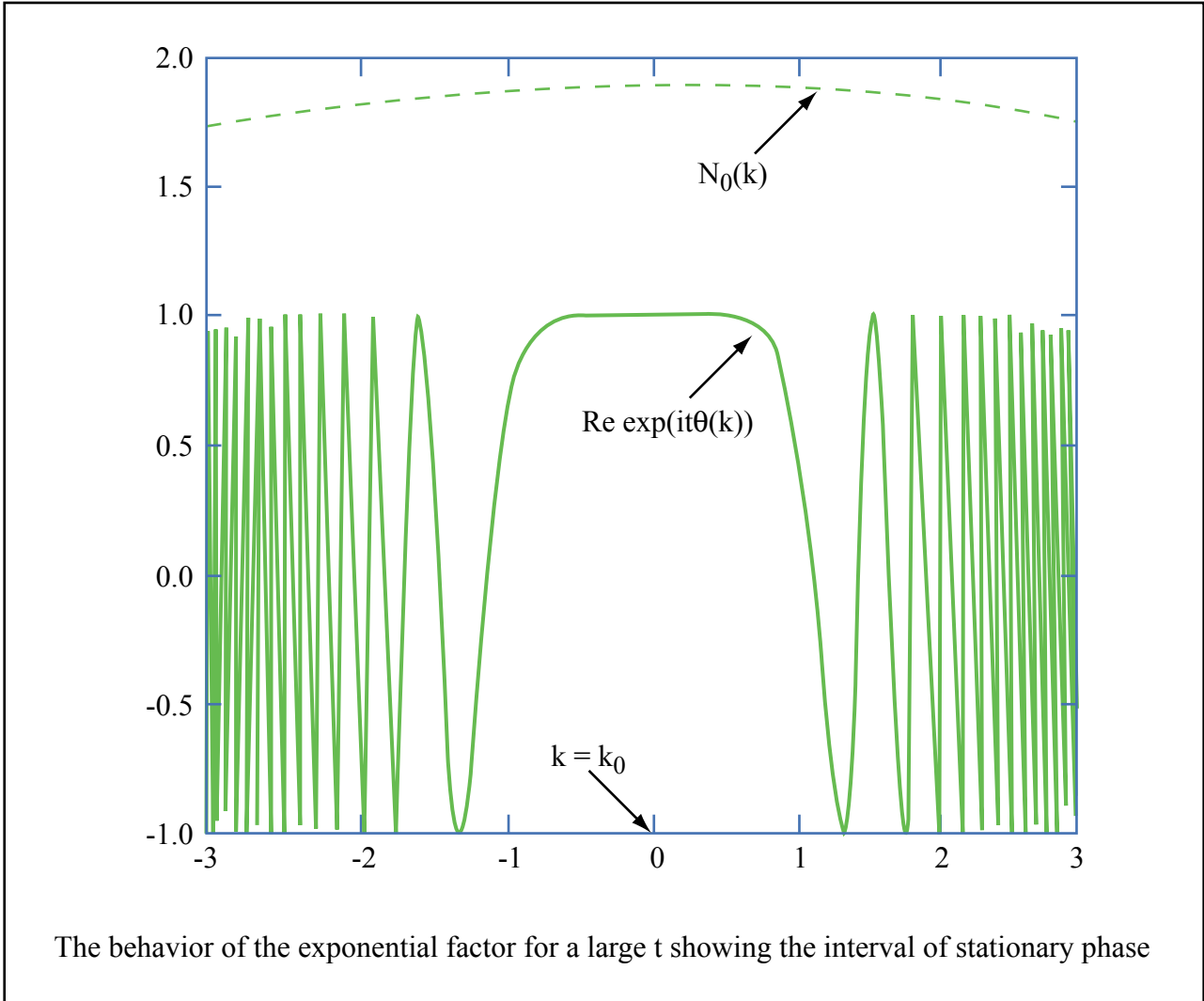
If we put a wavestaff at fixed  $x$  and record  $\eta(x, t)$  as a function of time



The wavelength decreases and the amplitude increases linearly

$t = 0$  amplitude = 0

Figure by MIT OpenCourseWare.



The behavior of the exponential factor for a large  $t$  showing the interval of stationary phase

Figure by MIT OpenCourseWare.

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