Lecture 6

PN Junction and MOS Electrostatics(III)

Metal-Oxide-Semiconductor Structure

Outline

1. Introduction to MOS structure
2. Electrostatics of MOS in thermal equilibrium
3. Electrostatics of MOS with applied bias

Reading Assignment:
Howe and Sodini, Chapter 3, Sections 3.7-3.8

1. Introduction

Metal-Oxide-Semiconductor structure

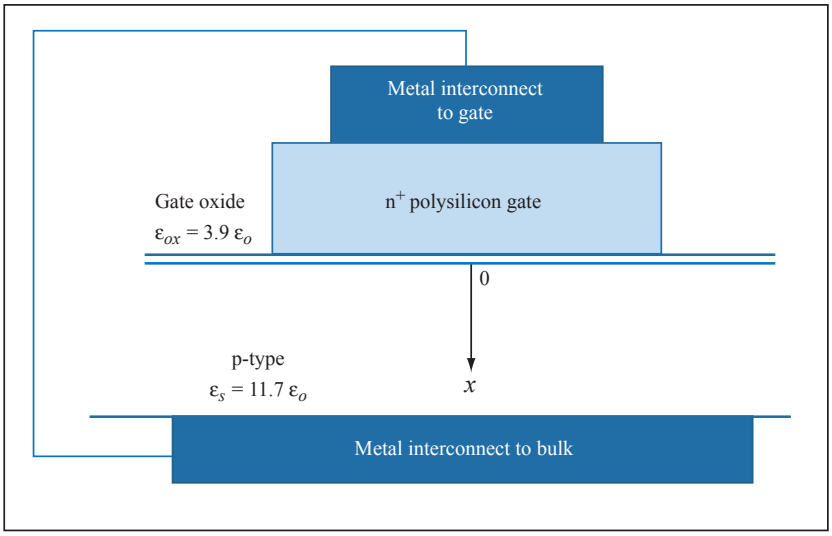


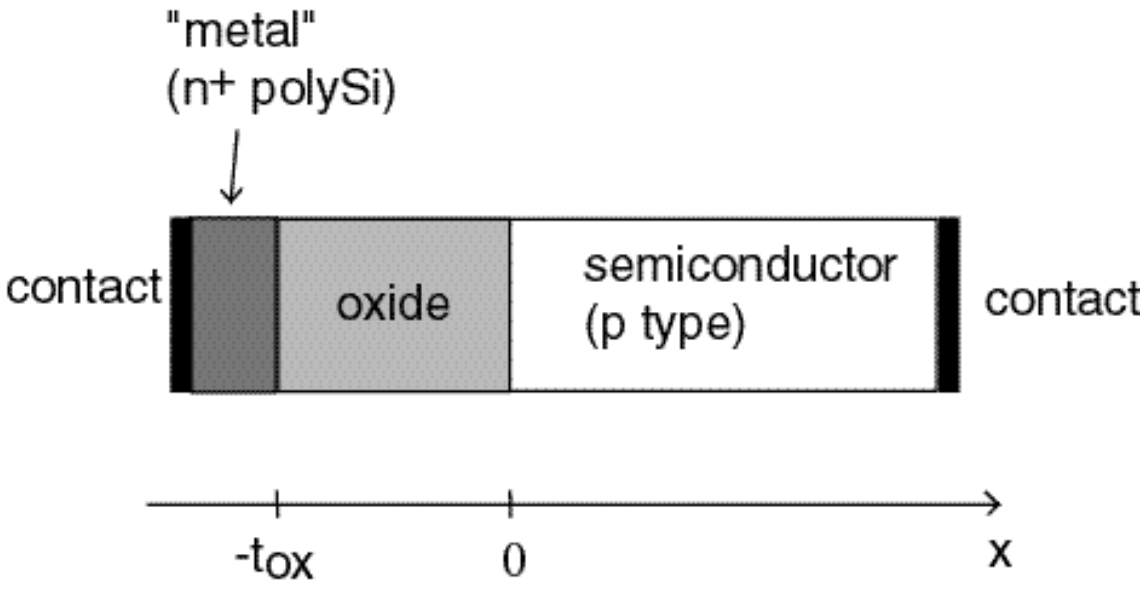
Figure by MIT OpenCourseWare.

MOS at the heart of the electronics revolution:

- Digital and analog functions
  - Metal-Oxide-Semiconductor Field-Effect Transistor (MOSFET) is key element of Complementary Metal-Oxide-Semiconductor (CMOS) circuit family
- Memory function
  - Dynamic Random Access Memory (DRAM)
  - Static Random Access Memory (SRAM)
  - Non-Volatile Random Access Memory (NVRAM)
- Imaging
  - Charge Coupled Device (CCD) and CMOS cameras
- Displays
  - Active Matrix Liquid Crystal Displays (AMLCD)

2. MOS Electrostatics in equilibrium

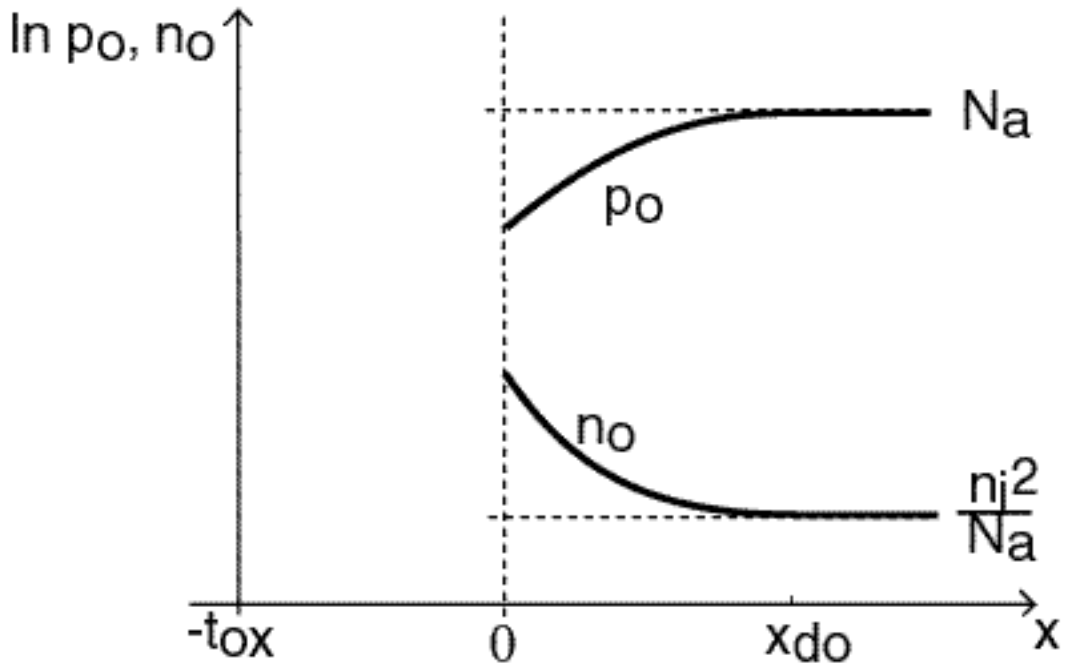
Idealized 1D structure:



- **Metal**: does not tolerate volume charge
  - ⇒ charge can only exist at its surface
- **Oxide**: insulator and does not have volume charge
  - ⇒ no free carriers, no dopants
- **Semiconductor**: can have volume charge
  - ⇒ Space charge region (SCR)

In thermal equilibrium we assume Gate contact is shorted to Bulk contact. (i. e, $V_{GB} = 0V$)

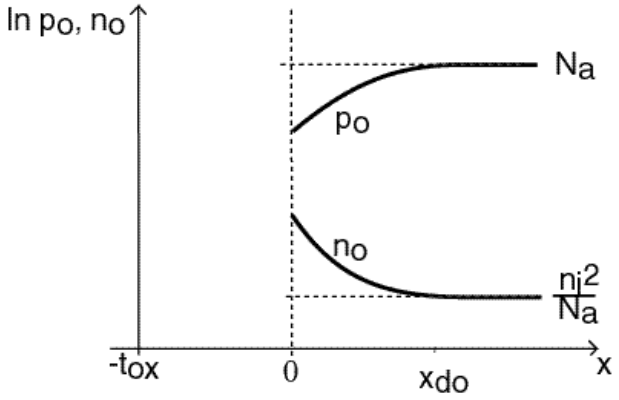
For most metals on p-Si, equilibrium achieved by electrons flowing from metal to semiconductor and holes from semiconductor to metal:

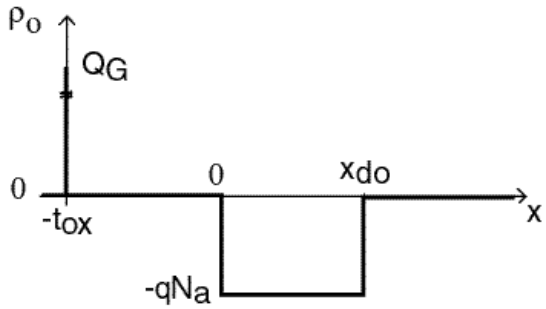


Remember: $n_o p_o = n_i^2$

Fewer holes near Si / SiO_2 interface
⇒ ionized acceptors exposed *(volume charge)*

Space Charge Density





- In semiconductor: space-charge region close Si /SiO_2 interface
  - can use ***depletion approximation***
- In metal: sheet of charge at metal /SiO_2 interface
- Overall charge neutrality

$$x = -t_{ox}; \qquad \sigma = Q_G$$

$$-t_{ox} < x < 0; \qquad \rho_o(x) = 0$$

$$0 < x < x_{do}; \qquad \rho_o(x) = -qN_a$$

$$x > x_{do}; \qquad \rho_o(x) = 0$$

Electric Field

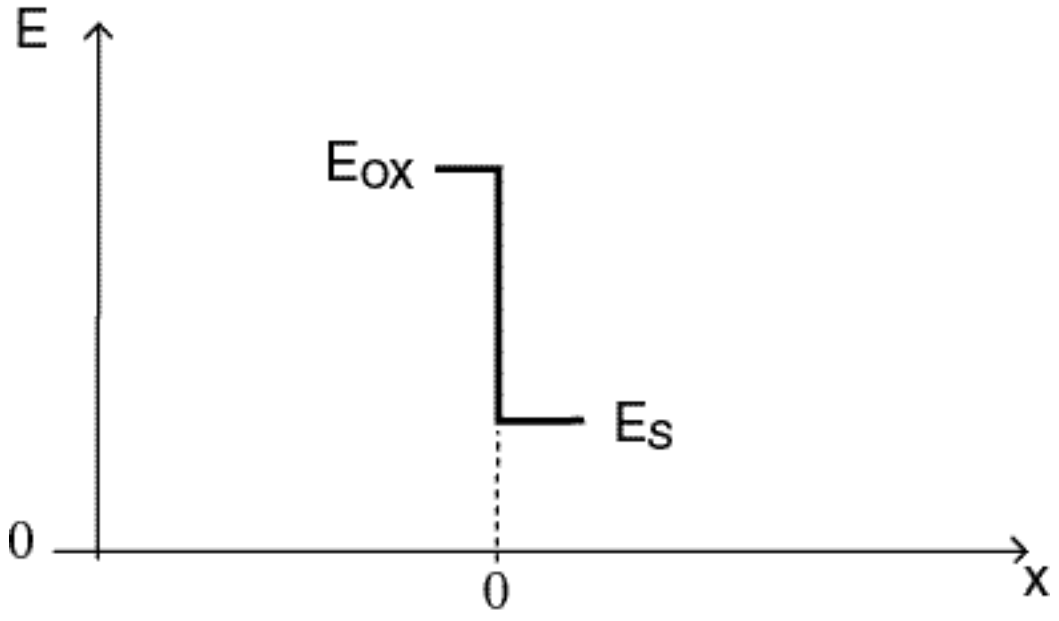
Integrate Poisson's equation

$$E_o(x_2) - E_o(x_1) = \frac{1}{\varepsilon}\int_{x_1}^{x_2} \rho(x')\,dx'$$

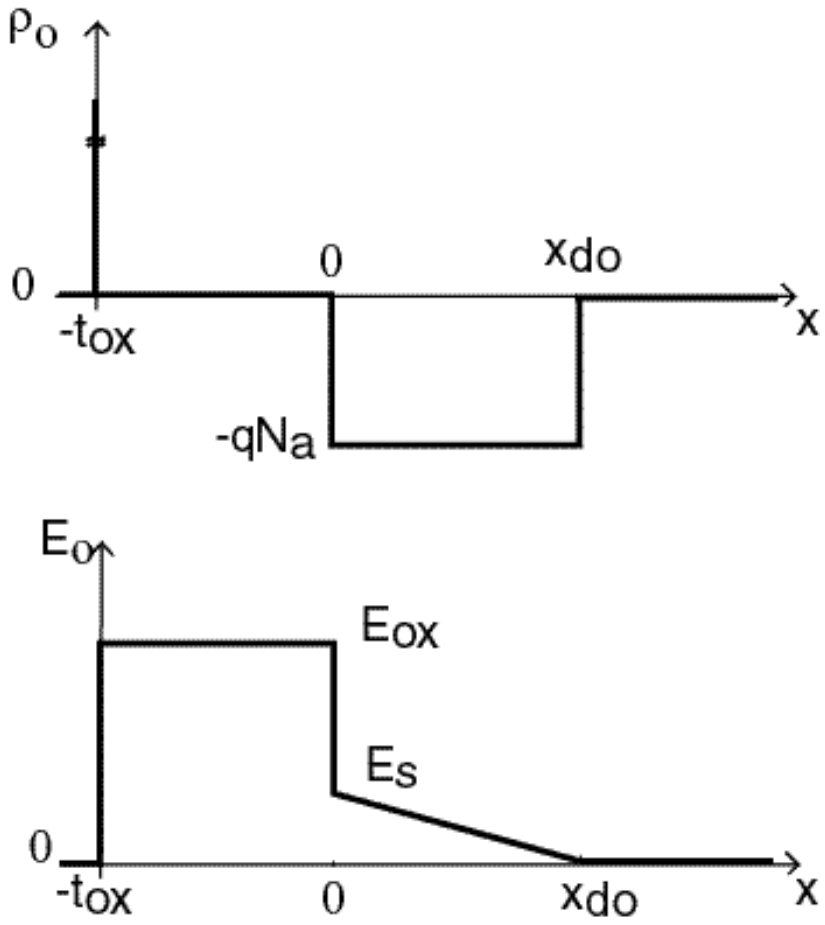
At interface between oxide and semiconductor, there is a change in **permittivity** $\Rightarrow$ change in electric field

$$\varepsilon_{ox} E_{ox} = \varepsilon_s E_s$$

$$\frac{E_{ox}}{E_s} = \frac{\varepsilon_s}{\varepsilon_{ox}} \approx 3$$



Start integrating from deep inside semiconductor:



$x > x_{do}$; $\quad E_o(x) = 0$

$0 < x < x_{do}$; $\quad E_o(x) - E_o(x_{do}) = \frac{1}{\varepsilon_s}\int_{x_{do}}^{x} -qN_a\, dx' = -\frac{qN_a}{\varepsilon_s}(x - x_{do})$

$-t_{ox} < x < 0$; $\quad E_o(x) = \frac{\varepsilon_s}{\varepsilon_{ox}} E_o(x = 0^+) = \frac{qN_a x_{do}}{\varepsilon_{ox}}$

$x < -t_{ox}$; $\quad E(x) = 0$

Electrostatic Potential

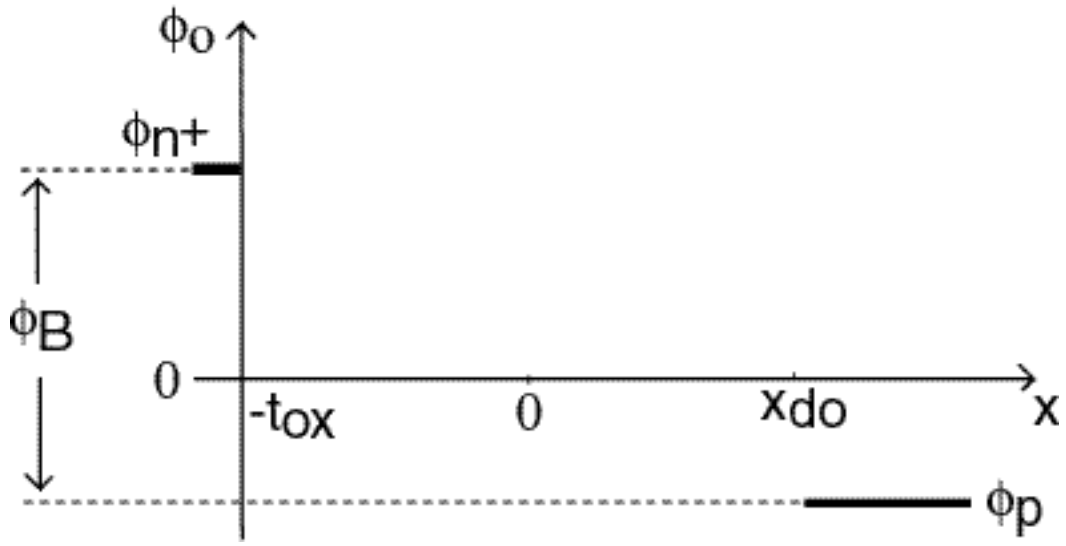
(with $\phi = 0$ @ $n_o = p_o = n_i$)

$$\phi = \frac{kT}{q} \bullet \ln \frac{n_o}{n_i} \qquad \phi = -\frac{kT}{q} \bullet \ln \frac{p_o}{n_i}$$

In QNRs, n_o and p_o are known $\Rightarrow$ can determine ϕ

in p-QNR: $p_o = N_a \Rightarrow \quad \phi_p = -\frac{kT}{q} \bullet \ln \frac{N_a}{n_i}$

in n^+-gate: $n_o = N_d^+ \Rightarrow \quad \phi_g = \phi_{n^+}$

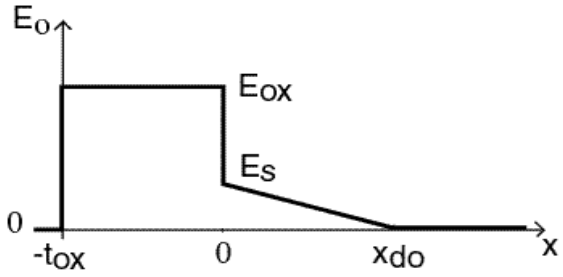


Built-in potential:

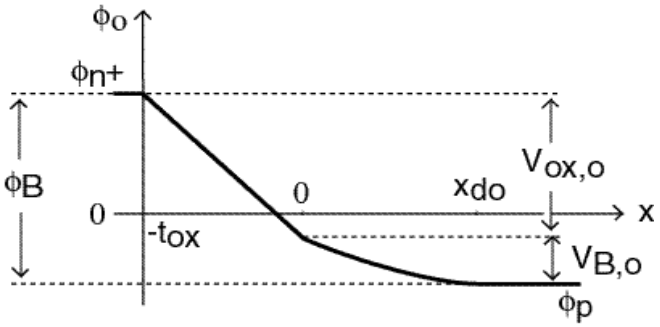
$$\phi_B = \phi_g - \phi_p = \phi_{n^+} + \frac{kT}{q} \bullet \ln \frac{N_a}{n_i}$$

To obtain $\phi_o(x)$, integrate $E_o(x)$; start from deep inside semiconductor bulk:

$$\phi_o(x_2) - \phi_o(x_1) = -\int_{x_1}^{x_2} E_o(x')\,dx'$$



$$\phi_B = \phi_{n^+} - \phi_p$$



$x > x_{do}$; $\quad \phi_o(x) = \phi_p$

$0 < x < x_{do}$; $\quad \phi_o(x) - \phi_o(x_{do}) = -\int_{x_{do}}^{x} -\frac{qN_a}{\varepsilon_s}(x' - x_{do})\,dx'$

$$\phi_o(x) = \phi_p + \frac{qN_a}{2\varepsilon_s}(x - x_{do})^2$$

$AT \quad x = 0 \qquad \phi_o(0) = \phi_p + \frac{qN_a}{2\varepsilon_s}(x_{do})^2$

$-t_{ox} < x < 0$; $\quad \phi_o(x) = \phi_p + \frac{qN_a x_{do}^2}{2\varepsilon_s} + \frac{qN_a x_{do}}{\varepsilon_{ox}}(-x)$

$x < -t_{ox}$; $\quad \phi_o(x) = \phi_{n^+}$

Almost done ….

Still do not know x_{do} $\Rightarrow$ need one more equation

Potential difference across structure has to add up to ϕ_B:

$$\phi_B = V_{B,o} + V_{ox,o} = \frac{qN_a x_{do}^2}{2\varepsilon_s} + \frac{qN_a x_{do} t_{ox}}{\varepsilon_{ox}}$$

Solve quadratic equation:

$$x_{do} = \frac{\varepsilon_s}{\varepsilon_{ox}} t_{ox} \left[\sqrt{1 + \frac{2\varepsilon_{ox}^2 \phi_B}{q\varepsilon_s N_a t_{ox}^2}} - 1 \right]$$

$$= \frac{\varepsilon_s}{C_{ox}} \left[\sqrt{1 + \frac{2C_{ox}^2 \phi_B}{q\varepsilon_s N_a}} - 1 \right]$$

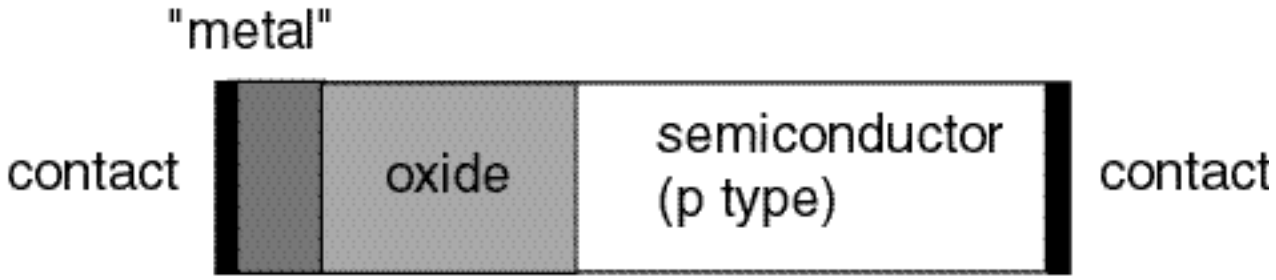
where C_{ox} is the capacitance per unit area of oxide

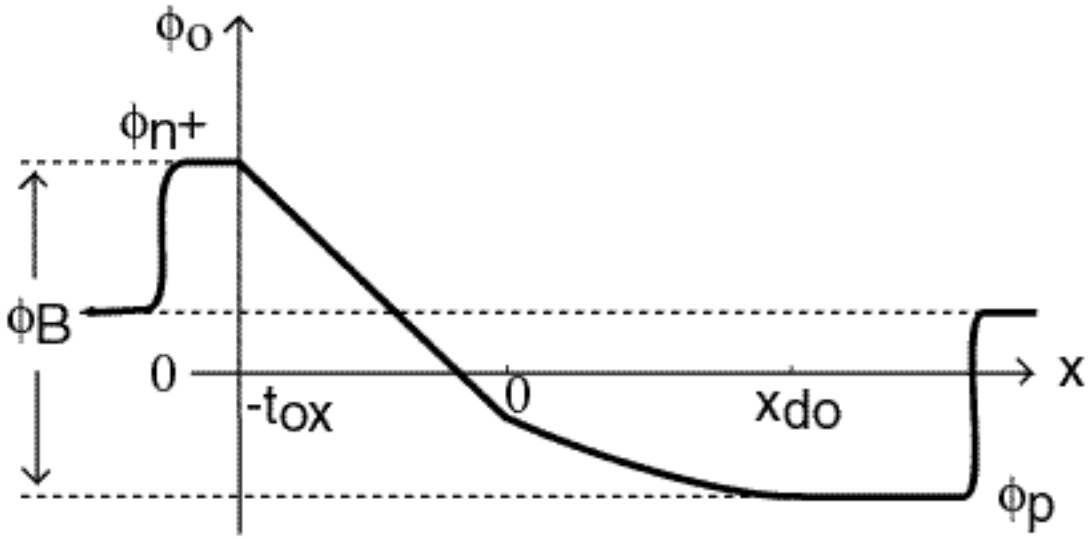
$$C_{ox} = \frac{\varepsilon_{ox}}{t_{ox}}$$

Now problem is completely solved!

There are also contact potentials

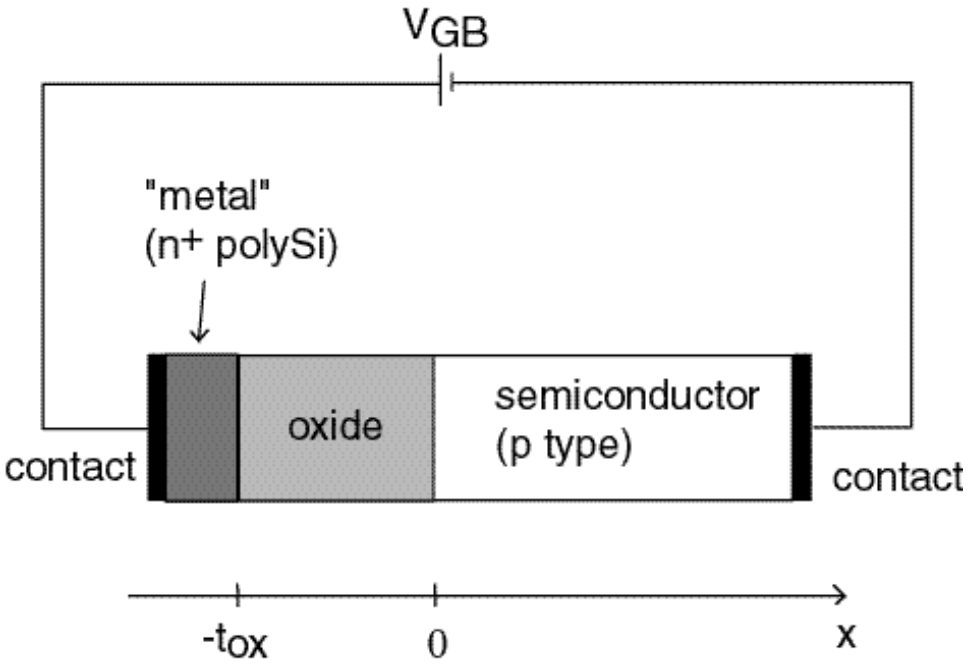
⇒ total potential difference from contact to contact is zero!





3. MOS with applied bias V_{GB}

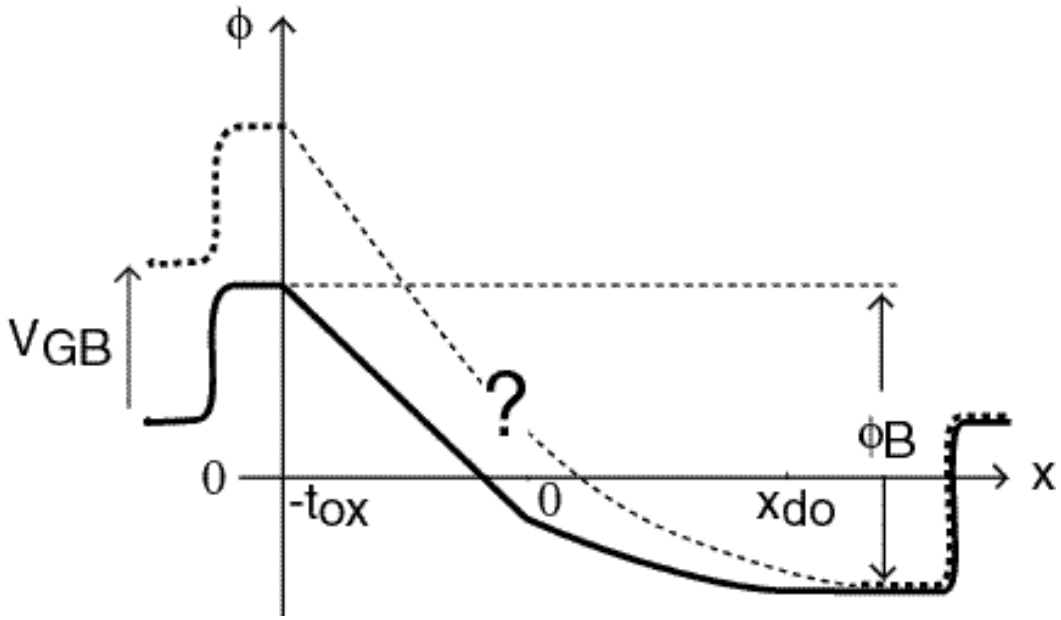
Apply voltage to gate with respect to semiconductor:



Electrostatics of MOS structure affected
$\Rightarrow$ potential difference across entire structure now $\neq 0$

How is potential difference accommodated?

Potential difference shows up across oxide and SCR in semiconductor

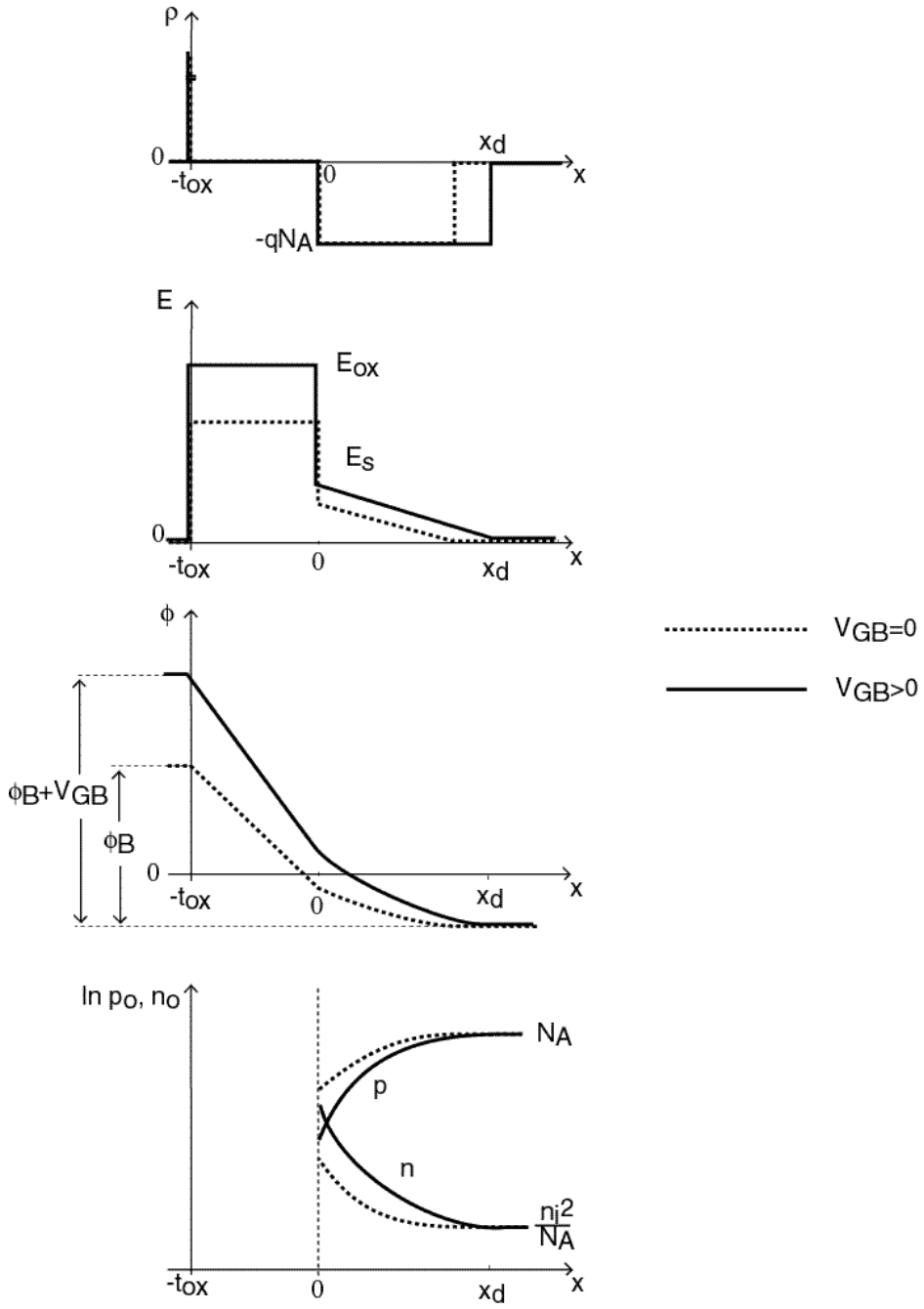


Oxide is an insulator ⇒ no current anywhere in structure

In SCR, quasi-equilibrium situation prevails
⇒ New balance between drift and diffusion

- Electrostatics qualitatively identical to thermal equilibrium (*but amount of charge redistribution is different*)
- $\mathbf{np = n_i^2}$

Apply $V_{GB}>0$: potential difference across structure increases ⇒ need larger charge dipole ⇒ SCR expands into semiconductor substrate:



Simple way to remember:
with $V_{GB}>0$, gate attracts electrons and repels holes.

Qualitatively, physics unaffected by application of $V_{GB} > 0$. Use mathematical formulation in thermal equilibrium, but:

$$\phi_B \rightarrow \phi_B + V_{GB}$$

For example, to determine $x_d(V_{BG})$:

$$\phi_B + V_{GB} = V_B(V_{GB}) + V_{ox}(V_{GB})$$

$$= \frac{qN_a x_d^2(V_{GB})}{2\varepsilon_s} + \frac{qN_a x_d(V_{GB}) t_{ox}}{\varepsilon_{ox}}$$

$$x_d(V_{GB}) = \frac{\varepsilon_s}{C_{ox}} \left[\sqrt{1 + \frac{2C_{ox}^2 (\phi_B + V_{GB})}{\varepsilon_s q N_a}} - 1 \right]$$

$$\phi(0) = \phi_s = \phi_p + \frac{qN_a x_d^2(V_{GB})}{2\varepsilon_s}$$

ϕ_s gives n & p concentration at the surface

What did we learn today?

Summary of Key Concepts

- Charge redistribution in MOS structure in thermal equilibrium
  - SCR in semiconductor
  - $\Rightarrow$ **built-in potential across MOS structure**.
- In most cases, we can use depletion approximation in semiconductor SCR
- Application of voltage modulates depletion region width in semiconductor
  - **No current flows**

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