

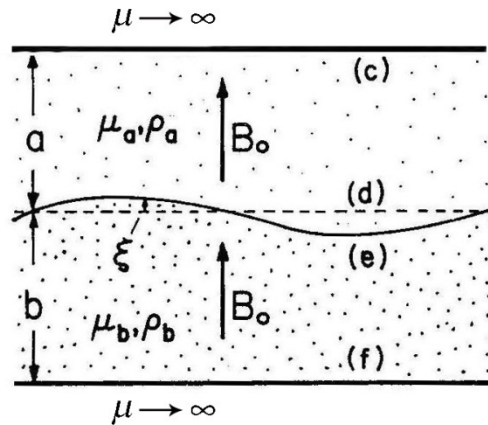
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6.642 Continuum Electromechanics
Fall 2008

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Lecture 8: Electrohydrodynamic and Ferrohydrodynamic instabilities

I. Magnetic Field Normal Instability



(a) Layers of magnetizable fluid are stressed by a uniform normal magnetic flux density, B_0 .

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A. Equilibrium ($\xi = 0$)

$$B_0 = \mu_a H_a = \mu_b H_b \quad ; \quad P_{oe} - P_{od} + \frac{1}{2} [\mu_a H_a^2 - \mu_b H_b^2] = 0$$

$$P_o(x) = \begin{cases} -\rho_a g x + P_{od} & x > 0 \\ -\rho_b g x + P_{oe} & x < 0 \end{cases}$$

B. Perturbations: $e^{j(\omega t - k_y y - k_z z)}$

$$\bar{H}_a = H_a \bar{i}_x + \bar{h}_a, \quad \bar{H}_b = H_b \bar{i}_x + \bar{h}_b$$

$$\begin{bmatrix} \hat{h}_{xc} \\ \hat{h}_{xd} \end{bmatrix} = k \begin{bmatrix} -\coth ka & \frac{1}{\sinh ka} \\ -\frac{1}{\sinh ka} & \coth ka \end{bmatrix} \begin{bmatrix} \hat{\Psi}^c \\ \hat{\Psi}^d \end{bmatrix}$$

$$\begin{bmatrix} \hat{h}_{xe} \\ \hat{h}_{xf} \end{bmatrix} = k \begin{bmatrix} -\coth kb & \frac{1}{\sinh kb} \\ -\frac{1}{\sinh kb} & \coth kb \end{bmatrix} \begin{bmatrix} \hat{\Psi}^e \\ \hat{\Psi}^f \end{bmatrix}$$

C. Boundary Conditions

$$\hat{v}_{xc} = \hat{v}_{xf} = 0$$

$$\hat{\psi}^c = \hat{\psi}^f = 0$$

$$v_{xd,e} = \frac{\partial \xi}{\partial t} + v_{yd,e} \frac{\partial \xi}{\partial y} + v_{zd,e} \frac{\partial \xi}{\partial z} \Rightarrow \hat{v}_{xd} = \hat{v}_{xe} = j\omega \hat{\xi}$$

$$\bar{n} = \frac{\bar{i}_x - \frac{\partial \xi}{\partial y} \bar{i}_y - \frac{\partial \xi}{\partial z} \bar{i}_z}{\left[1 + \left(\frac{\partial \xi}{\partial y} \right)^2 + \left(\frac{\partial \xi}{\partial z} \right)^2 \right]} \approx \bar{i}_x - \frac{\partial \xi}{\partial y} \bar{i}_y - \frac{\partial \xi}{\partial z} \bar{i}_z$$

$$\|p\| \mathbf{n}_i = \|T_{ij}^e\| \mathbf{n}_j - \gamma \nabla \cdot \bar{n} \mathbf{n}_i$$

$$\nabla \cdot \bar{n} = - \left(\frac{\partial^2 \xi}{\partial y^2} + \frac{\partial^2 \xi}{\partial z^2} \right)$$

$$\mathbf{i} = \mathbf{x} \Rightarrow \mathbf{n}_i = \mathbf{n}_x = 1$$

$$\|P_0 + P'\| = \|T_{xx}\| \underbrace{\|n_x\|}_{1} + \underbrace{\|T_{xy}\| \|n_y\| + \|T_{xz}\| \|n_z\|}_{\substack{\text{perturbation} \\ \mu H_x h_y \quad \mu H_x h_z}} + \gamma \left(\frac{\partial^2 \xi}{\partial y^2} + \frac{\partial^2 \xi}{\partial z^2} \right)$$

second order

Equilibrium ($\xi = 0$)

$$\|P_0\| = \|T_{xx0}\| \Rightarrow P_{od} - P_{oe} = \frac{1}{2} [\mu_a H_a^2 - \mu_b H_b^2]$$

Perturbations

$$P'_d(\xi) - P'_e(\xi) = \frac{1}{2} [2\mu_a H_a h'_d(\xi) - 2\mu_b H_b h'_e(\xi)] - \gamma k^2 \hat{\xi}$$

$$P'_d(\xi) = P_{od}(\xi) + P'_d(0) = \frac{dP_{od}}{dx} \Big|_{x=0} \xi + P'_d(0) = -\rho_a g \xi + P'_d(0)$$

$$P'_e(\xi) = P_{oe}(\xi) + P'_e(0) = \frac{dP_{oe}}{dx} \Big|_{x=0} \xi + P'_e(0) = -\rho_b g \xi + P'_e(0)$$

$$\begin{bmatrix} \hat{p}_c \\ \hat{p}_d \end{bmatrix} = \frac{j\omega\rho_a}{k} \begin{bmatrix} -\coth ka & \frac{1}{\sinh ka} \\ -\frac{1}{\sinh ka} & \coth ka \end{bmatrix} \begin{bmatrix} \hat{v}_x^c \\ \hat{v}_x^d \end{bmatrix}$$

$$\begin{bmatrix} \hat{p}_e \\ \hat{p}_f \end{bmatrix} = \frac{j\omega\rho_b}{k} \begin{bmatrix} -\coth kb & \frac{1}{\sinh kb} \\ -\frac{1}{\sinh kb} & \coth kb \end{bmatrix} \begin{bmatrix} \hat{v}_x^e \\ \hat{v}_x^f \end{bmatrix}$$

$$\hat{p}_d = \frac{j\omega\rho_a}{k} \coth ka \hat{v}_x^d = \frac{-\omega^2\rho_a}{k} \coth ka \hat{\xi}$$

$$\hat{p}_e = -\frac{j\omega\rho_b}{k} \coth kb \hat{v}_x^e = \frac{\omega^2\rho_b}{k} \coth kb \hat{\xi}$$

$$\hat{h}_{xd} = k \coth ka \hat{\Psi}^d$$

$$\hat{h}_{xe} = -k \coth kb \hat{\Psi}^e$$

$$\bar{n} \cdot \|\bar{b}\| = 0 = \left[\bar{i}_x - \frac{\partial \xi}{\partial y} \bar{i}_y - \frac{\partial \xi}{\partial z} \bar{i}_z \right] \cdot \left[(B_a - B_b) \bar{i}_x + (\mu_a h_{xd} - \mu_b h_{xe}) \bar{i}_x + (\mu_a h_{yd} - \mu_b h_{ye}) \bar{i}_y + (\mu_a h_{zd} - \mu_b h_{ze}) \bar{i}_z \right]$$

$$B_a = B_b = B_0$$

$$\mu_a h_{xd} = \mu_b h_{xe}$$

$$\bar{n} \times \|\bar{h}\| = 0 = \left[\bar{i}_x - \frac{\partial \xi}{\partial y} \bar{i}_y - \frac{\partial \xi}{\partial z} \bar{i}_z \right] \times \left[(H_a - H_b) \bar{i}_x + (h_{xd} - h_{xe}) \bar{i}_x + (h_{yd} - h_{ye}) \bar{i}_y + (h_{zd} - h_{ze}) \bar{i}_z \right]$$

$$= \bar{i}_z (h_{yd} - h_{ye}) - \bar{i}_y (h_{zd} - h_{ze}) + \frac{\partial \xi}{\partial y} (H_a - H_b) \bar{i}_z - \frac{\partial \xi}{\partial z} (H_a - H_b) \bar{i}_y$$

$$(h_{yd} - h_{ye}) = -\frac{\partial \xi}{\partial y} (H_a - H_b) \Rightarrow +jk_y (\hat{\Psi}_d - \hat{\Psi}_e) = jk_y (H_a - H_b) \hat{\xi}$$

$$(h_{zd} - h_{ze}) = -\frac{\partial \xi}{\partial z} (H_a - H_b) \Rightarrow +jk_z (\hat{\Psi}_d - \hat{\Psi}_e) = jk_z (H_a - H_b) \hat{\xi}$$

$$\hat{\Psi}_d - \hat{\Psi}_e = + (H_a - H_b) \hat{\xi}$$

$$\mu_a \mathcal{K} \coth ka \hat{\Psi}_d = -\mu_b \mathcal{K} \coth kb \hat{\Psi}_e$$

$$\hat{\Psi}_d = \frac{-\mu_b}{\mu_a} \frac{\coth kb}{\coth ka} \hat{\Psi}_e$$

$$\hat{\Psi}_e \left[-\frac{\mu_b}{\mu_a} \frac{\coth kb}{\coth ka} - 1 \right] = (H_a - H_b) \hat{\xi}$$

$$\hat{\Psi}_e = -\frac{(H_a - H_b) \hat{\xi}}{\mu_a \coth ka + \mu_b \coth kb} \mu_a \coth ka$$

$$\hat{\Psi}_d = +\frac{(H_a - H_b) \hat{\xi}}{\mu_a \coth ka + \mu_b \coth kb} \mu_b \coth kb$$

D. Dispersion Equation

$$\begin{aligned} & -\rho_a g \hat{\xi} - \frac{\omega^2 \rho_a}{k} \coth ka \hat{\xi} + \rho_b g \hat{\xi} - \frac{\omega^2 \rho_b}{k} \coth kb \hat{\xi} \\ & = \frac{\mu_a H_a k \coth ka (H_a - H_b) \mu_b \coth kb \hat{\xi} - \mu_b H_b k \coth kb (H_a - H_b) \hat{\xi} \mu_a \coth ka}{\mu_a \coth ka + \mu_b \coth kb} - \gamma k^2 \hat{\xi} \end{aligned}$$

$$\frac{\omega^2}{k} (\rho_a \coth ka + \rho_b \coth kb) = (\rho_b - \rho_a) g + \gamma k^2 - \frac{B_0 (H_a - H_b) \coth ka \coth kb (\mu_b - \mu_a) k}{\mu_a \coth ka + \mu_b \coth kb}$$

$$H_a = \frac{B_0}{\mu_a}, H_b = \frac{B_0}{\mu_b} \Rightarrow H_a - H_b = B_0 \left(\frac{1}{\mu_a} - \frac{1}{\mu_b} \right) = \frac{(\mu_b - \mu_a) B_0}{\mu_a \mu_b}$$

$$\frac{\omega^2}{k} (\rho_a \coth ka + \rho_b \coth kb) = g(\rho_b - \rho_a) + \gamma k^2 - \frac{B_0^2 (\mu_b - \mu_a)^2 k}{\mu_a \mu_b [\mu_b \tanh ka + \mu_a \tanh kb]}$$

E. Short Wavelength Limit ($ka \gg 1, kb \gg 1$)

$$\tanh ka \approx \tanh kb \approx 1$$

$$\frac{\omega^2}{k} (\rho_a + \rho_b) = g(\rho_b - \rho_a) + \gamma k^2 - \frac{B_0^2 (\mu_b - \mu_a)^2 k}{\mu_a \mu_b (\mu_a + \mu_b)} = f$$

Incipience of Instability

$$f = 0$$

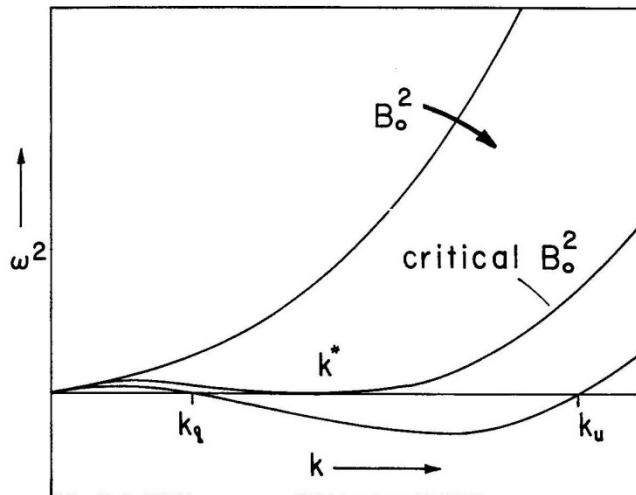
$$\frac{df}{dk} = 0 = 2\gamma k_c - \frac{B_0^2 (\mu_b - \mu_a)^2}{\mu_a \mu_b (\mu_a + \mu_b)}$$

$$k_c = \frac{B_0^2 (\mu_b - \mu_a)^2}{2\gamma \mu_a \mu_b (\mu_a + \mu_b)}$$

$$g(\rho_b - \rho_a) + \gamma \left[\frac{B_0^2 (\mu_b - \mu_a)^2}{2\gamma \mu_a \mu_b (\mu_a + \mu_b)} \right]^2 - \left[\frac{B_0^2 (\mu_b - \mu_a)^2}{\mu_a \mu_b (\mu_a + \mu_b)} \right]^2 \frac{1}{2\gamma} = 0$$

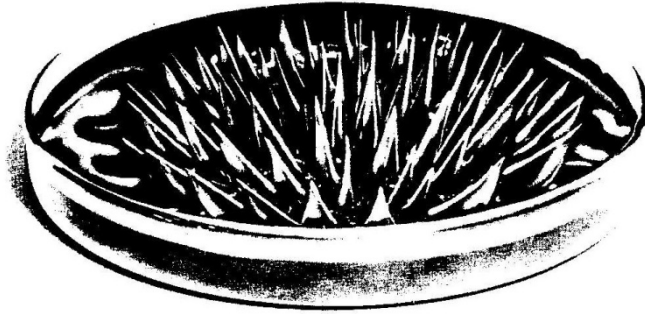
$$\left[\frac{B_0^2 (\mu_b - \mu_a)^2}{\mu_a \mu_b (\mu_a + \mu_b)} \right]^2 = 4g(\rho_b - \rho_a)\gamma$$

$$k_c = \frac{1}{2\gamma} \sqrt{4g(\rho_b - \rho_a)\gamma} = \sqrt{\frac{g(\rho_b - \rho_a)}{\gamma}}$$



Dependence of $\omega^2(k)$ as given by Eq. 18 with B_0^2 as a parameter.

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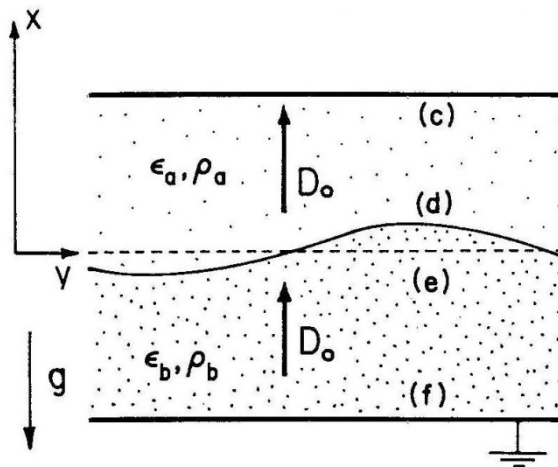


System of static fluid sprouts represents a new static equilibrium formed once planar interface in perpendicular field becomes unstable. (Courtesy of Ferrofluidics Corp., Burlington, Mass.)

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II. Electric Field Normal Instability

A. Polarization Forces



Polarizable liquid layers are stressed by a normal electric displacement, D_0 .

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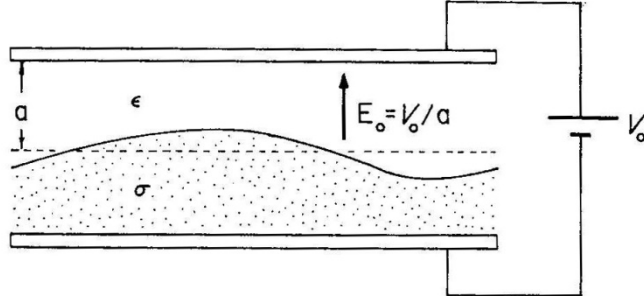
$$\mu_a \rightarrow \epsilon_a$$

$$\mu_b \rightarrow \epsilon_b$$

$$B_0 \rightarrow D_0$$

$$\frac{\omega^2}{k} (\rho_a \coth ka + \rho_b \coth kb) = g(\rho_b - \rho_a) + \gamma k^2 - \frac{D_0^2 (\epsilon_b - \epsilon_a)^2 k}{\epsilon_a \epsilon_b [\epsilon_b \tanh ka + \epsilon_a \tanh kb]}$$

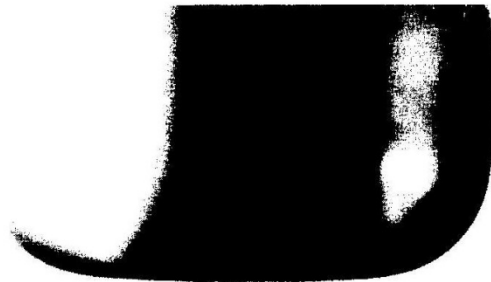
B. Perfectly conducting lower fluid ($\epsilon_b \rightarrow \infty$)



Rigid plane-parallel electrodes bound liquids having common interface. The upper liquid is insulating relative to the lower one.

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$$\frac{\omega^2}{k} (\rho_a \coth ka + \rho_b \coth kb) = g(\rho_b - \rho_a) + \gamma k^2 - \epsilon_a \left(\frac{V_0}{a}\right)^2 k \coth ka$$

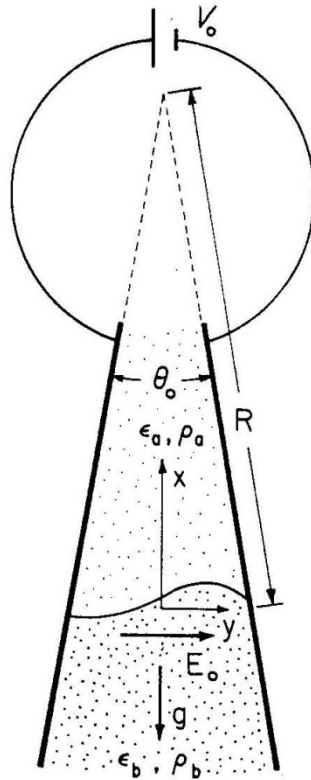


Nonlinear stages of surface instability caused by applying 30 kV d-c between electrode above and glycerine interface below.

Insulation is mixture of air and gaseous Freon.

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III. Tangential Gradient Fields



Cross section experiment
with Cartesian coordinates
for planar model.

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A. Equilibrium

$$\bar{E} = \bar{i}_\theta \frac{V_0}{\theta_0 r} \approx \bar{i}_y \frac{V_0}{\theta_0 r} \left(1 + \frac{x}{R}\right); r \approx R - x; E_0 = \frac{V_0}{\theta_0 R}$$

$$\approx \bar{i}_y E_0 \left(1 + \frac{x}{R}\right)$$

$$\lim_{\Delta \rightarrow \infty} P_{oa}(x) = -\rho_a g x + P_{oa}$$

$$P_{ob}(x) = -\rho_b g x + P_{ob}$$

$$P_{ob} - P_{oa} + \|T_{xx}\| = 0$$

$$T_{xx} = -\frac{1}{2} \varepsilon E_y^2 = -\frac{1}{2} \varepsilon E_0^2$$

$$\|T_{xx}\| = -\frac{1}{2} (\varepsilon_a - \varepsilon_b) E_0^2 = P_{oa} - P_{ob}$$

B. Perturbations

$$\begin{bmatrix} \hat{e}_x^\alpha \\ \hat{e}_x^\beta \end{bmatrix} = k \begin{bmatrix} -\coth k\Delta & \frac{1}{\sinh k\Delta} \\ -\frac{1}{\sinh k\Delta} & \coth k\Delta \end{bmatrix} \begin{bmatrix} \hat{\Phi}^\alpha \\ \hat{\Phi}^\beta \end{bmatrix}$$

$$\lim_{\Delta \rightarrow \infty}$$

$$\hat{e}_{xa} = k \hat{\Phi}_a$$

$$\hat{e}_{xb} = -k \hat{\Phi}_b$$

$$\begin{bmatrix} \hat{p}^\alpha \\ \hat{p}^\beta \end{bmatrix} = \frac{j\omega\rho}{k} \begin{bmatrix} -\coth k\Delta & \frac{1}{\sinh k\Delta} \\ -\frac{1}{\sinh k\Delta} & \coth k\Delta \end{bmatrix} \begin{bmatrix} \hat{v}_x^\alpha \\ \hat{v}_x^\beta \end{bmatrix}$$

$$\hat{v}_x^\alpha = \hat{v}_x^\beta = j\omega \hat{\xi}$$

$$\hat{p}_a = \frac{j\omega\rho_a}{k} \hat{v}_{xa} = -\frac{\omega^2 \rho_a}{k} \hat{\xi}$$

$$\hat{p}_b = -\frac{j\omega\rho_b}{k} \hat{v}_{xb} = \frac{\omega^2\rho_b}{k} \hat{\xi}$$

C. Boundary Conditions

$$\bar{n} = \bar{i}_x - \frac{\partial \xi}{\partial y} \bar{i}_y - \frac{\partial \xi}{\partial z} \bar{i}_z$$

$$\bar{n} \times \|\bar{E}\| = 0 \Rightarrow \left[\bar{i}_x - \frac{\partial \xi}{\partial y} \bar{i}_y - \frac{\partial \xi}{\partial z} \bar{i}_z \right] \times \left[(e_{xa} - e_{xb}) \bar{i}_x + (e_{ya} - e_{yb}) \bar{i}_y + (e_{za} - e_{zb}) \bar{i}_z \right] = 0$$

$$\bar{i}_z (e_{ya} - e_{yb}) - \bar{i}_y (e_{za} - e_{zb}) = 0$$

$$\left. \begin{aligned} e_{ya} = e_{yb} &\Rightarrow jk_y \hat{\Phi}_a = jk_y \hat{\Phi}_b \\ e_{za} = e_{zb} &\Rightarrow jk_z \hat{\Phi}_a = jk_z \hat{\Phi}_b \end{aligned} \right\} \Rightarrow \hat{\Phi}_a = \hat{\Phi}_b \equiv \hat{\Phi}$$

$$\bar{n} \cdot \|\varepsilon \bar{E}\| = 0 \Rightarrow \left[\bar{i}_x - \frac{\partial \xi}{\partial y} \bar{i}_y - \frac{\partial \xi}{\partial z} \bar{i}_z \right] \cdot \left[E_0 (\varepsilon_a - \varepsilon_b) \bar{i}_y + (\varepsilon_a e_{xa} - \varepsilon_b e_{xb}) \bar{i}_x + (\varepsilon_a e_{ya} - \varepsilon_b e_{yb}) \bar{i}_y + (\varepsilon_a e_{za} - \varepsilon_b e_{zb}) \bar{i}_z \right] = 0$$

$$\varepsilon_a \hat{e}_{xa} - \varepsilon_b \hat{e}_{xb} + jk_y \hat{\xi} E_0 (\varepsilon_a - \varepsilon_b) = 0$$

$$k \hat{\Phi} (\varepsilon_a + \varepsilon_b) + jk_y E_0 (\varepsilon_a - \varepsilon_b) \hat{\xi} = 0$$

$$\hat{\Phi} = \frac{-jk_y E_0 (\varepsilon_a - \varepsilon_b) \hat{\xi}}{k (\varepsilon_a + \varepsilon_b)}$$

$$\|\rho\| n_i = \|T_{ij}\| n_j + \gamma \left[\frac{\partial^2 \xi}{\partial y^2} + \frac{\partial^2 \xi}{\partial z^2} \right] n_i$$

$$i = x, n_x = 1$$

$$\|\hat{p}\| = \|\hat{T}_{xx}\| + \|\hat{T}_{xy}\| \hat{n}_y + \|\hat{T}_{xz}\| \hat{n}_z - \gamma k^2 \hat{\xi}$$

$$T_{xx} = \frac{1}{2} \varepsilon (E_x^2 - E_y^2 - E_z^2) = -\frac{1}{2} \varepsilon [E_{y0} + e_y]^2 = -\frac{1}{2} \varepsilon [E_{y0}^2 + 2E_0 e_y]$$

$$T_{xy} = \varepsilon E_{y0} e_x$$

$$T_{xz} = \varepsilon e_x e_z$$

$\|T_{xy}\|n_y$ second order

$\|T_{xz}\|n_z$ third order

$$\begin{aligned} T_{xx}(\xi) &= -\frac{1}{2} \varepsilon \frac{d}{dx} (E_{y0}^2) \Big|_{x=0} \xi - \varepsilon E_0 e_y \\ &= -\varepsilon E_0 \frac{dE_0}{dx} \xi - \varepsilon E_0 e_y \end{aligned}$$

$$\begin{aligned} \hat{T}_{xx} &= -\varepsilon E_0 \frac{dE_0}{dx} \hat{\xi} - \varepsilon E_0 \hat{e}_y \\ &= -\varepsilon E_0 \frac{dE_0}{dx} \hat{\xi} - jk_y \hat{\Phi} \varepsilon E_0 \end{aligned}$$

$$\|\hat{T}_{xx}\| = -(\varepsilon_a - \varepsilon_b) E_0 \frac{dE_0}{dx} \hat{\xi} - jk_y E_0 \hat{\Phi} (\varepsilon_a - \varepsilon_b)$$

$$\|\hat{p}_T\| = \left\| \frac{dp_0}{dx} \hat{\xi} + \hat{p} \right\| = \hat{p}_a - \hat{p}_b - g(\rho_a - \rho_b) \hat{\xi}$$

$$\hat{p}_a - \hat{p}_b - g(\rho_a - \rho_b) \hat{\xi} = -(\varepsilon_a - \varepsilon_b) E_0 \frac{dE_0}{dx} \hat{\xi} - jk_y \hat{\Phi} E_0 (\varepsilon_a - \varepsilon_b) - \gamma k^2 \hat{\xi}$$

$$-\frac{\omega^2}{k} (\rho_a + \rho_b) \hat{\xi} - g(\rho_a - \rho_b) \hat{\xi} =$$

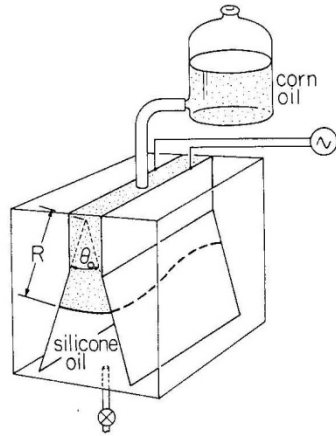
$$-(\varepsilon_a - \varepsilon_b) E_0 \frac{dE_0}{dx} \hat{\xi} - \gamma k^2 \hat{\xi} - \frac{jk_y E_0 (\varepsilon_a - \varepsilon_b) (-jk_y E_0) (\varepsilon_a - \varepsilon_b) \hat{\xi}}{k (\varepsilon_a + \varepsilon_b)}$$

$$\frac{\omega^2}{k} (\rho_a + \rho_b) = g(\rho_b - \rho_a) + \gamma k^2 + (\varepsilon_a - \varepsilon_b) E_0 \frac{dE_0}{dx} + \frac{k_y^2 E_0^2 (\varepsilon_a - \varepsilon_b)^2}{k (\varepsilon_a + \varepsilon_b)}$$

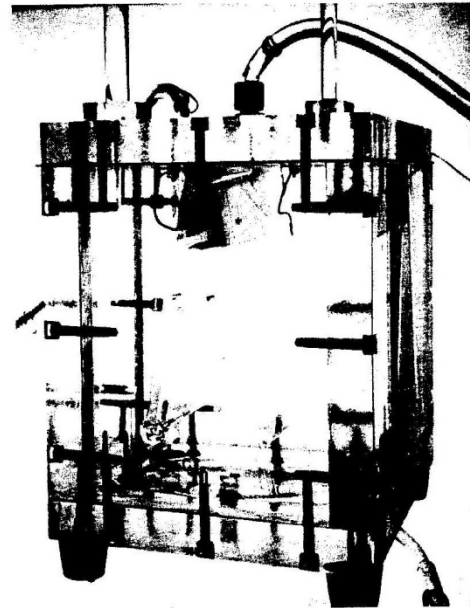
Uniform tangential field always stabilizes.

Gradient field stabilizes, if higher permittivity fluid is in stronger electric field. System stable even if heavier fluid above if

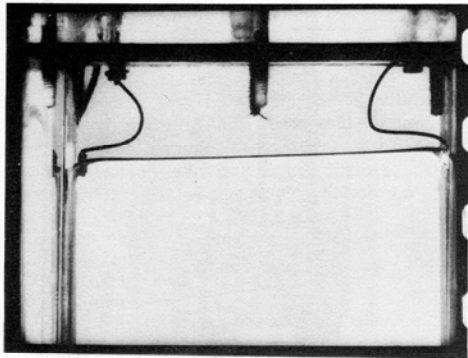
$$(\varepsilon_a - \varepsilon_b) E_0 \frac{dE_0}{dx} > g(\rho_a - \rho_b)$$



Heavy liquid is stabilized on top of lighter fluid by means of polarization forces induced by applying potential difference to the diverging glass plates. These plates have a thin transparent coating that renders them conducting.

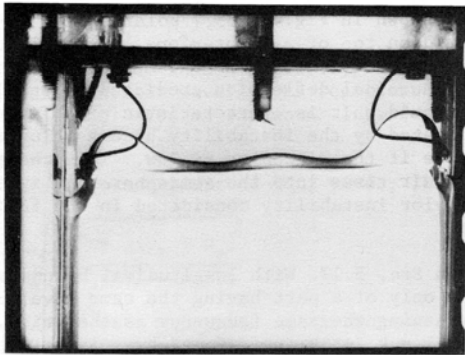


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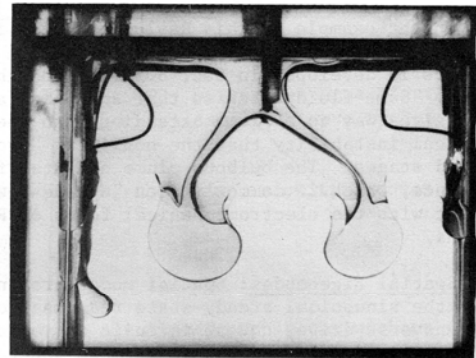


(a)

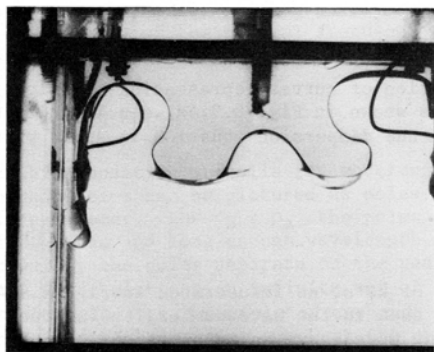
Side view of apparatus.
(a) Equilibrium with field on.
(b)-(e) Sequential view of developing instability.
(From Complex Waves II, Reference 11, Appendix C.)



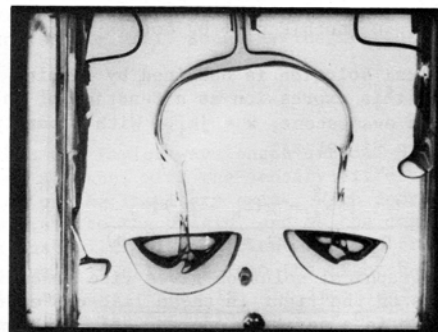
(b)



(d)



(c)



(e)

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