

Perfect Conductivity Lecture 2

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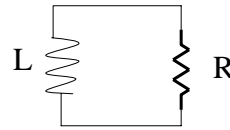
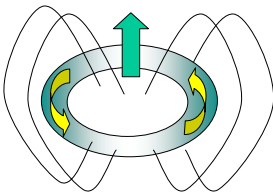
Outline

1. **Persistent Currents**
2. **Parts of a Physical Theory**
3. **Circuits and Time Constants**
4. **Distributive Systems and Time constants**
 - A. **Quasistatics**
 - B. **MagnetoQuasiStatics (MQS)**

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Persistent Currents



If the field is turned off, then

$$I(t) = I|_{t=0} e^{-t/\tau_{LR}}$$

The time constant $\tau_{LR} = L/R$

If the loop is made out of a superconductor,

$$\lim_{R \rightarrow 0} \tau_{LR} \rightarrow \infty \quad I(t) = I|_{t=0} \quad \text{for } t \geq 0.$$

Experimentally the dc resistivity of a superconductor is at least as small as $10^{-25} \Omega\text{-m}$. The superconducting state is “truly” zero dc resistance.

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Perfect Conductivity: $t \ll \tau_{RL} = L/R$,
system looks like R is zero

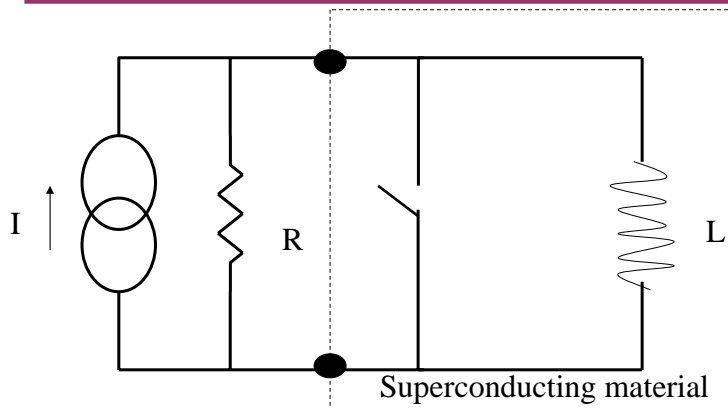
Superconductivity: for all time,
R is zero

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Charging up a superconducting loop



This Persistent Mode is the basis of MRI magnets, SMES, flux memory....



Parts of a Physical Theory

1. Governing Laws:

Maxwell's Equations, Newton's equations,

2. Constitutive Laws:

Models of the system

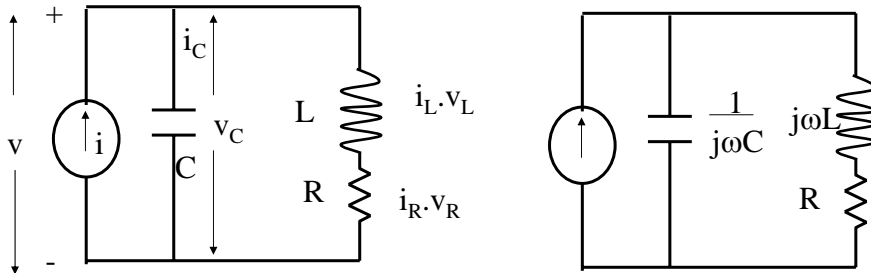
like ohm's law,

3. Summary Relations:

Transfer functions, Dispersion relations



LRC Circuit



1. Governing Equations

$$\text{Current conservation: } i = i_C + i_L \quad i_L = i_R$$

$$\text{Energy Conservation } v = v_C = v_R + v_L$$



2. Constitutive Relations

For the resistor

$$v_R = i_R R,$$

for the inductor as

$$v_L = L \frac{d}{dt} i_L,$$

and for the capacitor as

$$i_C = C \frac{d}{dt} v_C,$$

and

so that

$$\hat{i} \equiv |i| e^{j\phi}$$

$$v = \text{Re} \{ \hat{v} e^{j\omega t} \}$$

$$v = \text{Re} \{ \hat{v} e^{j\omega t} \}$$

$$\hat{v}_R = \hat{i}_R R,$$

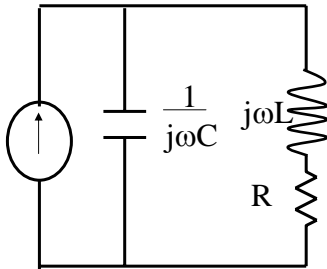
$$\hat{v}_L = j\omega L \hat{i}_L$$

$$\hat{i}_C = j\omega C \hat{v}_C$$



3. Summary Relation

$$Z(\omega) = R \left(\frac{1 + j\omega\tau_{RL}}{(1 - (\omega\tau_{LC})^2) + j\omega\tau_{RC}} \right).$$



the inductive time constant:

$$\tau_{RL} \equiv \frac{L}{R}, \quad (1)$$

the capacitive time constant:

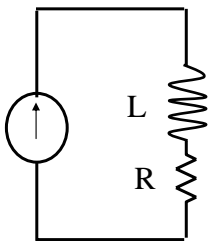
$$\tau_{RC} \equiv RC, \quad (2)$$

and the coupling time constant:

$$\tau_{LC} \equiv \sqrt{LC} = \sqrt{\tau_{RL}\tau_{RC}}. \quad (3)$$

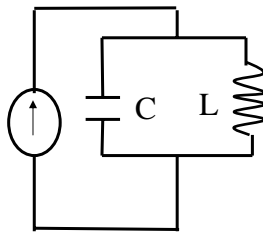


Simpler Circuits and Time Constants



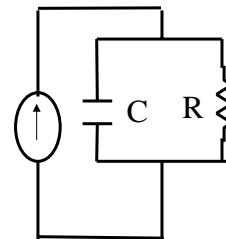
$$\tau_{RL} \equiv \frac{L}{R}$$

Energy stored
in inductor



$$\tau_{LC} \equiv \sqrt{LC}$$

Resonant transfer of
energy between L and C

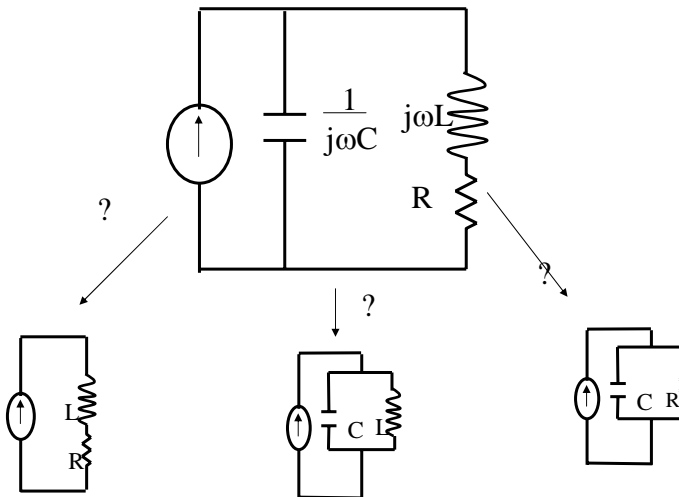


$$\tau_{RC} \equiv RC$$

Energy stored in
capacitor



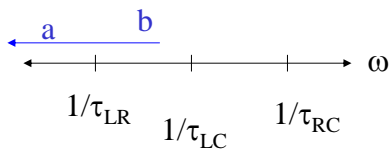
Reducing the Circuit to a simpler form ?



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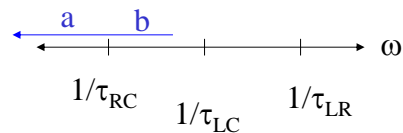
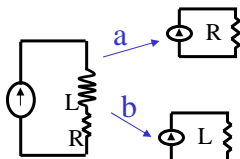
Order of time constants



$$\tau_{LR} > \tau_{RC} \quad \text{Low } R \quad R < \sqrt{\frac{L}{C}}$$

$$\lim_{\substack{\omega\tau_{LC} \ll 1 \\ \tau_{RC} < \tau_{RL}}} Z(\omega) \approx R(1 + j\omega\tau_{RL})$$

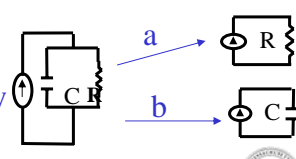
Low frequency circuit



$$\tau_{LR} < \tau_{RC} \quad \text{High } R \quad R > \sqrt{\frac{L}{C}}$$

$$\lim_{\substack{\omega\tau_{LC} \ll 1 \\ \tau_{RL} < \tau_{RC}}} Z(\omega) \approx R \left(\frac{1}{1 + j\omega\tau_{RC}} \right)$$

Low frequency circuit



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Moral of time constants

If you know what frequency range you want to study or what physics dominates the problem,
then you can solve a simpler problem.

Useful, especially in more complex situations.



Distributed Systems

1. Governing Equations: Maxwell's Equations

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \quad \text{Faraday's Law}$$

$$\nabla \times \mathbf{H} = \frac{\partial \mathbf{D}}{\partial t} + \mathbf{J} \quad \text{Ampere's Law}$$

$$\nabla \cdot \mathbf{D} = \rho \quad \text{Gauss Law}$$

$$\nabla \cdot \mathbf{B} = 0 \quad \text{Gauss' Magnetic Law}$$

Conservations laws

$$\nabla \cdot \mathbf{J} + \frac{\partial \rho}{\partial t} = 0 \quad \text{Charge conservation} \quad \text{Also Poynting's}$$



Distributed systems con't

2. Constitutive Relations

$$\mathbf{B}(\mathbf{r}, \omega) = \mu(\omega) \mathbf{H}(\mathbf{r}, \omega)$$

$$\mathbf{D}(\mathbf{r}, \omega) = \epsilon(\omega) \mathbf{E}(\mathbf{r}, \omega) \quad \begin{array}{l} \text{Local in space,} \\ \text{linear time invariant} \end{array}$$

$$\mathbf{J}(\mathbf{r}, \omega) = \sigma(\omega) \mathbf{E}(\mathbf{r}, \omega) \longrightarrow \text{Ohm's Law}$$

3. Summary relations

Complex: Search first for first order in time approximation



Quasistatic Limit

$$\ell \ll \lambda_{em} = \frac{2\pi c}{\omega}$$

Length scale of system ℓ

Wavelength of E&M wave λ_{em}

Speed of light c

Frequency (angular) ω

If the dimensions of a structure are much less than the wavelength of an electromagnetic field interacting with it, the coupling between the associated electric and magnetic fields is weak and a quasistatic approximation is appropriate.



Time Constants

$$\tau_{em} \equiv \frac{\ell}{c} = \ell \sqrt{\mu\epsilon} \quad \text{Electromagnetic coupling time}$$

$$\tau_e \equiv \frac{\epsilon}{\sigma_0} \quad \text{Charge relaxation time}$$

$$\tau_m \equiv \mu\sigma_0\ell^2 \quad \text{Magnetic diffusion time}$$

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Please see: Figure 2.9, page 34, from Orlando, T., and K. Delin. *Foundations of Applied Superconductivity*. Reading, MA: Addison-Wesley, 1991. ISBN: 0201183234.

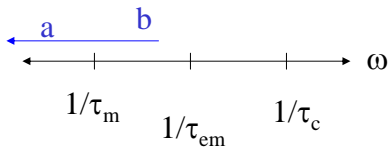
$$\tau_{em} = \sqrt{\tau_e\tau_m}$$

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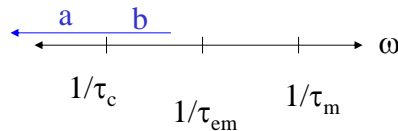
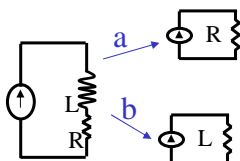
Order of time constants



$$\tau_m > \tau_c \quad \text{High conductivity} \quad \sigma_0 > \frac{1}{\ell} \sqrt{\frac{\epsilon}{\mu}}$$

MQS

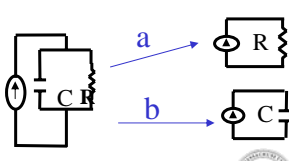
Low frequency circuit



$$\tau_m < \tau_e \quad \text{Low conductivity} \quad \sigma_0 < \frac{1}{\ell} \sqrt{\frac{\epsilon}{\mu}}$$

EQS

Low frequency circuit



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MagnetoQuasiStatics

$$\left. \begin{aligned} \nabla \times \mathbf{H} &= \mathbf{J} \\ \nabla \cdot \mathbf{B} &= 0 \\ \nabla \cdot \mathbf{J} &= 0 \end{aligned} \right\} \text{Solve first}$$

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$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \quad \text{Solve for } \mathbf{E} \text{ once } \mathbf{B} \text{ is found}$$

Boundary conditions:

$$\mathbf{n} \times (\mathbf{H}_2 - \mathbf{H}_1) = \mathbf{K} \quad \mathbf{n} \cdot (\mathbf{B}_2 - \mathbf{B}_1) = 0 \quad \mathbf{n} \cdot (\mathbf{J}_2 - \mathbf{J}_1) = 0$$

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MQS: Magnetic Diffusion Equation

For a metal $\mathbf{B} = \mu_0 \mathbf{H}$, $\mathbf{D} = \epsilon_0 \mathbf{E}$ and $\mathbf{J} = \sigma_0 \mathbf{E}$, so that

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$$\left(\mu \sigma_0 \frac{\partial}{\partial t} - \nabla^2 \right) \mathbf{H} = 0$$

Magnetic Diffusion Equation

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