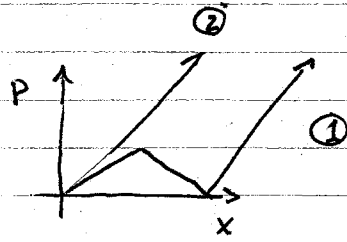


Th F
23
30
7
13 No class

Wed. 12 @ noon
Cancel Fri 7

Pick date for oral presentations

Recall confusion from last time

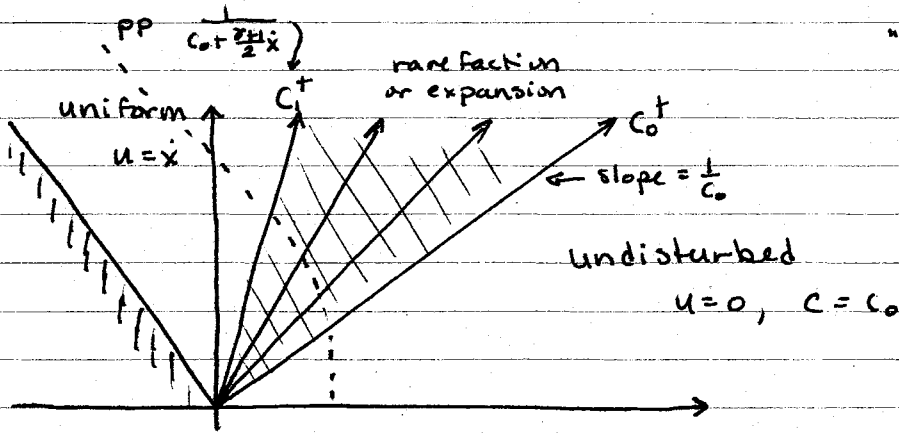
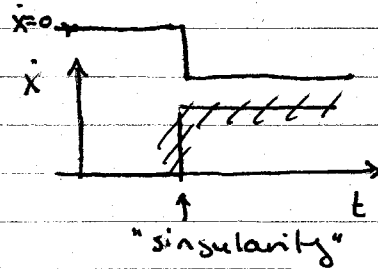


"The portion of a wave that increases the density as it passes is called a compression."

$[P] > 0$ compression shock $(P_2 - P_1) > 0$

Piston Withdrawal

$$\dot{x} = \begin{cases} 0 & t < 0 \\ \text{const} < 0 & t > 0 \end{cases}$$



singular pt.
multiple values for $x \therefore$
multiple values for u, c

Recall C^+ characteristics are straight lines

$$\text{slope} = \frac{1}{u+c}$$

\therefore slopes of characteristics vary between

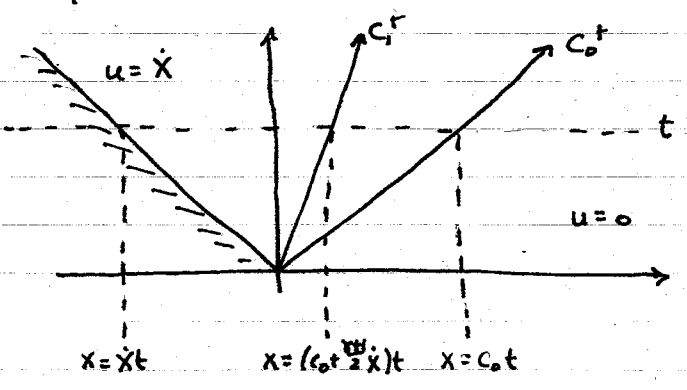
$$\frac{1}{u+c_0} = \frac{1}{c_0} \quad \text{and} \quad \frac{1}{\dot{x}+c_1}$$

Recall from last time (since $j^- = \text{const}$)

$$c = c_0 + \frac{\gamma-1}{2} u$$

$$\Rightarrow \frac{1}{x+ct} = \frac{1}{x+c_0 + \frac{\gamma-1}{2} x} = \frac{1}{c_0 + \frac{\gamma+1}{2} x}$$

This is all the info we need to find u and all thermodynamic quantities for all x and t .



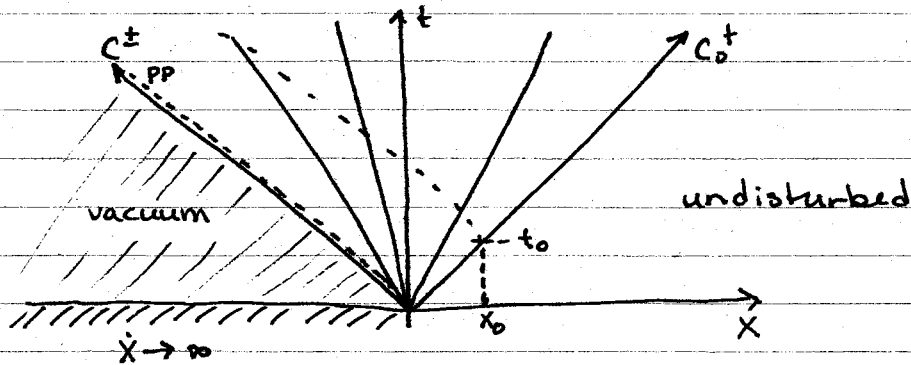
$\infty > x \geq c_0 t$	$\left\{ \begin{array}{l} u=0 \\ c=c_0 \end{array} \right.$	undisturbed
$c_0 t \geq x \geq (c_0 + \frac{\gamma+1}{2} X) t$	$\left\{ \begin{array}{l} u = \frac{2}{\gamma+1} (\frac{x}{t} - c_0) \\ c = c_0 + \frac{\gamma-1}{\gamma+1} (\frac{x}{t} - c_0) \end{array} \right.$	expansion wave
$(c_0 + \frac{\gamma+1}{2} X) t \geq x \geq X t$	$\left\{ \begin{array}{l} u = X \\ c = c_0 + \frac{\gamma-1}{2} X \end{array} \right.$	uniform

$$\frac{1}{\text{slope}} = \frac{x}{t} = u+c = u+c_0 + \frac{\gamma-1}{2} u \Rightarrow \frac{x}{t} = c_0 + \frac{\gamma+1}{2} u$$

$$\Rightarrow u = \frac{2}{\gamma+1} (\frac{x}{t} - c_0)$$

$$c = c_0 + \frac{\gamma-1}{2} \frac{2}{\gamma+1} (\frac{x}{t} - c_0)$$

This solution is ok provided piston moves slower than u_{escape} . If $|X| > |u_{\text{escape}}|$ then gas separates.



Recall $u_{\text{escape}} = -2c_0/(\gamma-1)$

~~$\frac{x}{t} = u + c = -2c_0/(\gamma-1)$~~

Solution as above except:

$x > c_0 t \quad \begin{cases} u=0 \\ c=c_0 \end{cases} \quad \text{undisturbed}$

$c_0 t > x > -2c_0/(\gamma-1) \quad \begin{cases} \text{same as expansion} \\ \text{wave above} \end{cases}$

$-2c_0/(\gamma-1) > x > -\infty \quad \text{vacuum} \quad (p=0, P=0)$

Particle path follows: $u = \frac{dx}{dt} = \frac{2}{\gamma+1} \left(\frac{x}{t} - c_0 \right)$ First order ODE for $x(t)$

$\Rightarrow x = -\frac{2}{\gamma-1} c_0 t + A t^{2/(\gamma+1)}$

↑ need to find this const. through b.c.

Apply b.c. @ $c_0 t_0 = x_0$

$$x_0 = -\frac{2}{\gamma-1} c_0 t_0 + A t_0^{2/(\gamma+1)}$$

$$x_0 + \frac{2}{\gamma-1} c_0 t_0 = A t_0^{2/(\gamma+1)} \Rightarrow A = t_0^{-2/(\gamma+1)} \left(x_0 + \frac{2}{\gamma-1} c_0 t_0 \right)$$

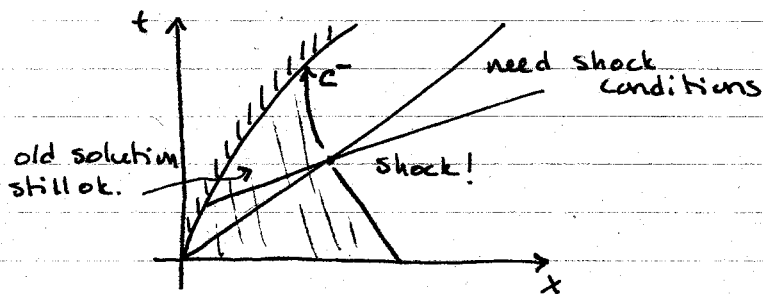
$$\Rightarrow x = -\frac{2}{\gamma-1} c_0 t + t_0^{-2/(\gamma+1)} \left(x_0 + \frac{2}{\gamma-1} c_0 t_0 \right) t^{2/(\gamma+1)}$$

$$= -\frac{2}{\gamma-1} c_0 t + \left(\frac{t}{t_0} \right)^{2/(\gamma+1)} x_0 \left(1 + \frac{2}{\gamma-1} \right)$$

$$x = -\frac{2}{\gamma-1} c_0 t + \frac{\gamma+1}{\gamma-1} x_0 \left(\frac{c_0 t}{x_0} \right)^{2/(\gamma+1)}$$

Continuous Piston Advance

Now $\dot{x} > 0$



Again characteristics end @ piston: $u+c = c_0 + \frac{\gamma+1}{2} \dot{x}$

\Rightarrow slope decreases as $t \uparrow$
if piston is accelerating

After shock forms, need ~~shock~~ shock conditions

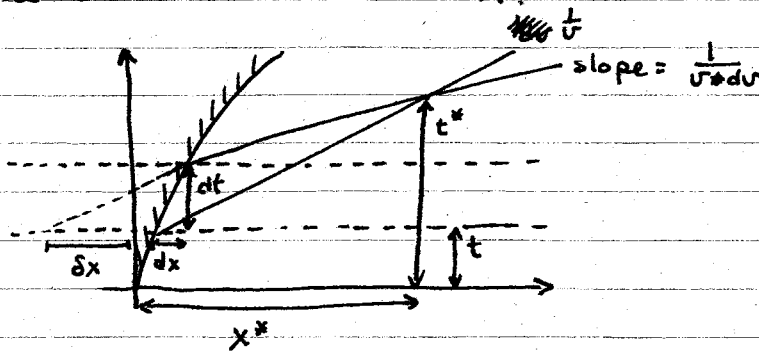
Solution plan: (1) Find if shock occurs

(2) Find when + where shock occurs

(3) Apply characteristic before shock

(4) Apply shock conditions after

Time of shock formation



$$v = \frac{x^*}{t^* - t}$$

$$v + dv = \frac{\delta x + dx}{dt} = \frac{\delta x + x^*}{t^* - t}$$

$$v + dv = \frac{\delta x}{t^* - t} + v$$

$$\Rightarrow t^* - t = d(utc)$$

Need δx and $d(utc)$

$$u + c = c_0 + \frac{\gamma + 1}{2} \dot{x} \Rightarrow d(utc) = \frac{\gamma + 1}{2} \ddot{x} dt$$

$$dt(v + \overset{\text{small}}{dv}) = \delta x + dx$$

$$\delta x = v dt - dx = v dt - \dot{x} dt$$

$$= (c_0 + \frac{\gamma + 1}{2} \dot{x} - \dot{x}) dt$$

$$= (c_0 + \frac{\gamma - 1}{2} \dot{x}) dt$$

$$\therefore t^* - t = \frac{(c_0 + \frac{\gamma - 1}{2} \dot{x}) dt}{\frac{\gamma + 1}{2} \ddot{x} dt}$$

$$t^* = t + \frac{1}{\dot{x}} \left(\frac{2}{\gamma + 1} c_0 + \frac{\gamma - 1}{\gamma + 1} \dot{x} \right)$$

Minimum (a) $\frac{dt^*}{dt} = 0$

$$0 = 1 + \frac{\ddot{x} \left(\frac{\gamma-1}{\gamma+1} \ddot{x} \dot{x} \right) - \left(\frac{2}{\gamma+1} c_0 + \frac{\gamma-1}{\gamma+1} \dot{x} \right) \ddot{x}}{\ddot{x}^2}$$

$$1 = \frac{\left(\frac{2}{\gamma+1} c_0 + \frac{\gamma-1}{\gamma+1} \dot{x} \right) \ddot{x}}{\ddot{x}^2} - \frac{\gamma-1}{\gamma+1}$$

$$\frac{2\gamma}{\gamma+1} \ddot{x}^2 = \left(\frac{2}{\gamma+1} c_0 + \frac{\gamma-1}{\gamma+1} \dot{x} \right) \ddot{x}$$

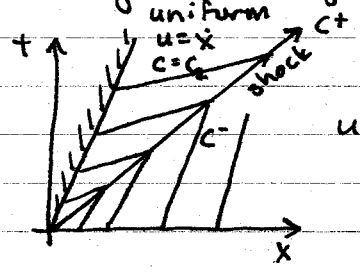
$$\boxed{\ddot{x}^2 = \frac{1}{\gamma} \left(c_0 + \frac{\gamma-1}{2} \dot{x} \right) \ddot{x}}$$

Implicit eq. for t_{min}^*

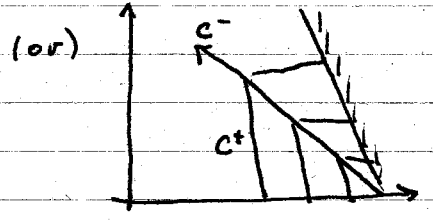
(Note: $x(t)$ is a ~~known~~ given fn of t)

Impulsively started piston

As before, singularity (shock this time) forms immediately @ origin



undisturbed
 $u=0, c=c,$



Step (4) Weak Shocks

Recall $[\rho] \sim [P]^3 \Rightarrow \approx$ isentropic

Assume J^- (or J^+) is const across shock (then go back and check if this is true.

$$J_1^- = J_2^- \quad J^- = u - \int dP/\rho c$$

$$u_2 - u_1 = \int_{P_1}^{P_2} \frac{dP}{\rho c} \quad \left(\frac{\partial u}{\partial P} \right)_s = - \frac{1}{\rho^2 c^2}$$

$$u_2 - u_1 = \int_{P_1}^{P_2} \sqrt{-\left(\frac{\partial u}{\partial P} \right)_s} dP \quad (*)$$

Taylor series:

$$\left(\frac{\partial u}{\partial P} \right)_s = \left(\frac{\partial u}{\partial P} \right)_s + \left(\frac{\partial^2 u}{\partial P^2} \right)_s (P - P_1) + \frac{1}{2} \left(\frac{\partial^3 u}{\partial P^3} \right)_s (P - P_1)^2 + \dots$$

\uparrow \uparrow
 $\Theta 2$ $\Theta 1$

Put this in (*) (approx $\sqrt{\quad}$) and integrate

$$\frac{u_2 - u_1}{c_1} = \frac{\pi}{c_1} - \frac{\frac{1}{2} \pi^2}{\frac{1}{2} \pi^2} - \frac{1}{6} \left[\pi^2 + \frac{c_1^6}{2u_1^4} \left(\frac{\partial^3 u}{\partial P^3} \right)_s \right] \pi^3 + \dots$$

\uparrow
 $\frac{[P]}{\rho_1 c_1}$

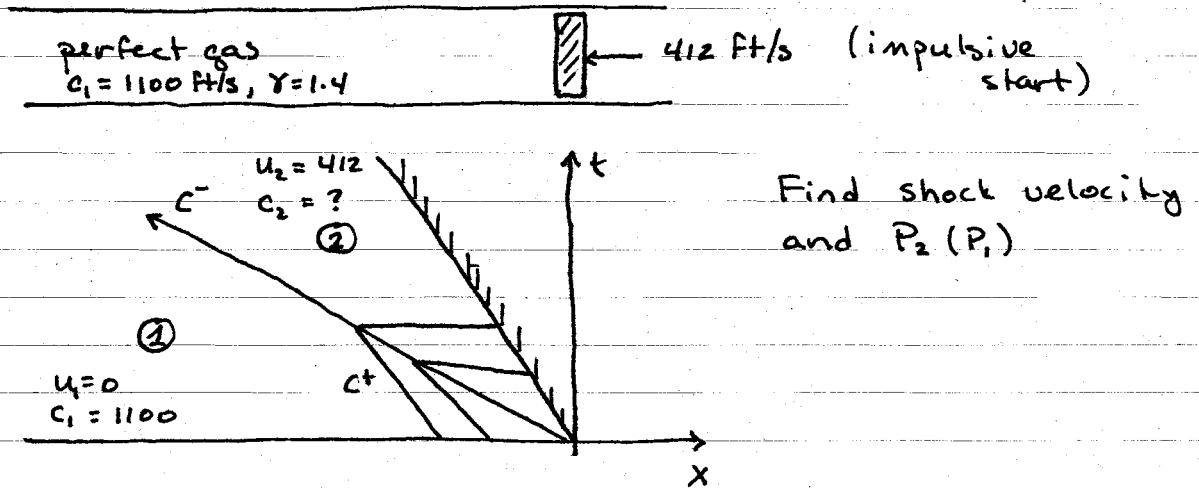
Using Taylor series, we can show that

$$[J] = \Theta(\pi^3) \quad (\text{same } \Theta \text{ as entropy!})$$

Treat as const across shock.)

Also follows that the shock velocity is approximately the average of the upstream and downstream wave velocities

$$V_{\text{shock}} = \frac{1}{2} [(u_1 + c_1) + (u_2 + c_2)] + \underbrace{c_1 \epsilon}_{\text{error } \Theta(\pi^2)}$$

Example

Assume shock is weak (and check later that this is not violated.)

$$[J^+] = 0$$

$$u_2 + \frac{\gamma}{\gamma-1} c_2 = u_1 + \frac{\gamma}{\gamma-1} c_1$$

$$\begin{aligned} & \text{or} \\ & -412 \text{ ft/s} \quad c_2 = 1100 \text{ ft/s} + 412 \text{ ft/s} \left(\frac{\gamma-1}{2} \right) = \boxed{1182 \text{ ft/s}} \end{aligned}$$

Shock vel. is (approximately) average of $(u-c)$ on both sides

$$V_{\text{shock}} = \frac{1}{2} [(-412 - 1182) + (0 - 1100)] = \boxed{-1347 \text{ ft/s}}$$

Isentropic perfect gas:

$$\frac{P_2}{P_1} = \left(\frac{v_1}{v_2} \right)^\gamma$$

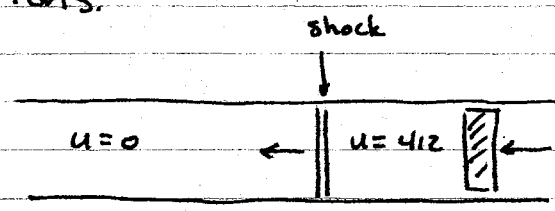
$$c^2 = \gamma P v$$

$$\frac{P_2}{P_1} = \left(\frac{c_1^2 \gamma P_1}{\gamma P_2 c_2^2} \right)^\gamma \Rightarrow \left(\frac{P_2}{P_1} \right)^{1-\gamma} = \left(\frac{c_1}{c_2} \right)^{2\gamma}$$

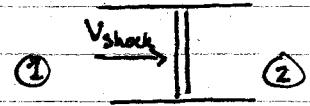
$$\Rightarrow \left(\frac{P_2}{P_1} \right) = \left(\frac{c_2}{c_1} \right)^{2\gamma/(\gamma-1)}$$

$$P_2 = P_1 \left(\frac{1182}{1100} \right)^{2.8/0.4} = \boxed{1.654 P_1}$$

But we could have also found this using normal shock relations.



Move to frame where shock is stationary



$$M_{in} = \frac{u_1}{c_1} = \frac{V_{shock}}{c_1}$$

$$-\frac{[u]}{c_1} = \frac{412}{1100} = 0.37455 \Rightarrow M_{in} = 1.25$$

↑
from tables

$$V = -1.25 (1100 \text{ ft/s}) = \boxed{-1375 \text{ ft/s}}$$

$$\frac{P_2}{P_1} = \boxed{1.656}$$

↑
from tables

Is weak approx ok? check that $\Pi \ll 1$

$$\Pi = \frac{[P]}{\delta P_1} = \frac{(1.656 - 1)P_1}{1.4 P_1} = 0.469 \quad (\text{approx. is borderline ok...})$$