

Solutions to Exam #3 8.03 Spring 2014

① a) $I = \frac{I_0}{2}$ $E = E_0 \sin(kz - \omega t) \hat{x} + E_0 \cos(kz - \omega t) \hat{y}$
 after polarizer only one \cos or \sin remains. e.g. $E_0 \sin(kz - \omega t) \hat{x}$

b) $\Delta t \cdot \Delta \nu = 1$ $\Delta \nu = 10^9 \text{ Hz} \sim 1 \text{ GHz}$

c) $v_p = \frac{\omega}{k} = \sqrt{\frac{g}{k}}$ $v_g = \frac{d\omega}{dk} = \frac{\sqrt{g}}{2\sqrt{k}}$ $v_p = 2v_g$

d) Width of the diffractive peak is $\approx \frac{\pi D}{\lambda} \sin \psi = \pi$
 $\sin \psi = \frac{R}{L} = \frac{\lambda}{D}$ $R = \frac{\lambda}{D} \cdot L = \frac{5 \cdot 10^{-7}}{1} \cdot 3.8 \cdot 10^8 = 200 \text{ m}$

More precise radius is obtained by analyzing circular sources $r = \frac{1.22 \lambda}{D} \cdot R = 230 \text{ m}$

e) Light reflected from water surface is predominantly polarized with \vec{E} parallel to water surface. Polaroid glasses are filtering out this polarization. By rotating by 90° we let them in, negating sun-blocking effect.

f) Use $\frac{\sin^2(\frac{N\delta}{2})}{\sin^2(\frac{\delta}{2})}$ $N \cdot d = \cos \theta \cdot t$ d smaller \Rightarrow Distance between principal maxima increases
 Width smaller, height larger
 λ larger Distance between PM increases

$$\textcircled{2} \quad a) \quad \vec{B} \perp \vec{k}, \quad \vec{\nabla} \cdot \vec{B} = 0 \quad \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} + \frac{\partial B_z}{\partial z} = 0$$

$$0 + \underbrace{\frac{\partial}{\partial y} (B_0 \cos(kz - \omega t))}_{=0!} + \underbrace{\omega \cdot \frac{\partial}{\partial z} (B_0 \cos(kz - \omega t))}_{\neq 0} = 0$$

$$\Rightarrow \omega = 0!$$

$$\vec{B}(\vec{r}, t) = B_0 \hat{y} \cos(kz - \omega t)$$

$$b) \quad \vec{E}(\vec{r}, t) = B_0 \cdot c \hat{x} \cos(kz - \omega t)$$

$$c) \quad \vec{F} = q(\vec{E} + \vec{v} \times \vec{B}) \quad \vec{a} = \frac{\vec{F}}{m}$$

Force due to \vec{B} field can be ignored since q is constrained to move along \hat{x}

$$\vec{a} = \frac{q \vec{E}}{m} = \frac{q B_0 c}{m} \hat{x} \cos(\omega t) \quad \text{since } z=0!$$

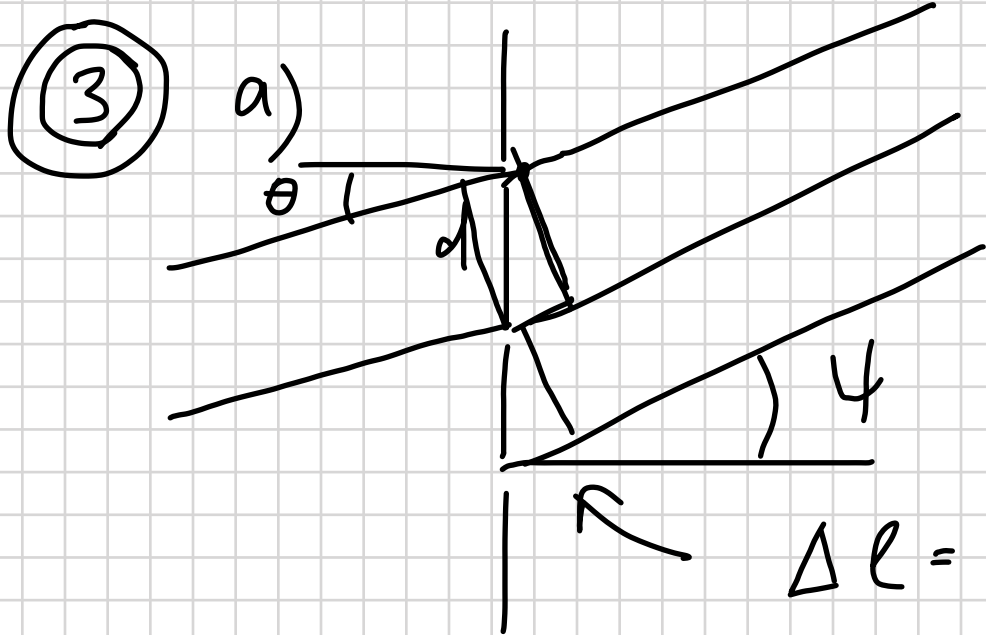
$$d) \quad \vec{E}_{r < d} = -\frac{q \vec{a}_\perp (t - rk)}{4\pi \epsilon_0 r c^2}$$

$$1) \quad \vec{r} = r \hat{x} \Rightarrow \vec{E}_{r < d} = 0, \quad \vec{a} \text{ is along } \hat{x}$$

$$2) \quad \vec{r} = r \hat{y} \Rightarrow \vec{E}_{r < d} = \frac{-q^2 B_0 c \cos(t - \frac{r}{c})}{4\pi \epsilon_0 r c^2} \hat{x}$$

Linear polarization

$$3) \quad \vec{r} = r \hat{z} \quad \text{SAME AS IN 2)}$$



$$\delta_{\text{after}} = \Delta l \cdot k = \frac{2\pi}{\lambda} \cdot d \sin \psi$$

b) Incoming wave will introduce add. phase delay. We defined θ and ψ such that phases cancel each other. In addition the presence of glass changes phase shift magnitude.

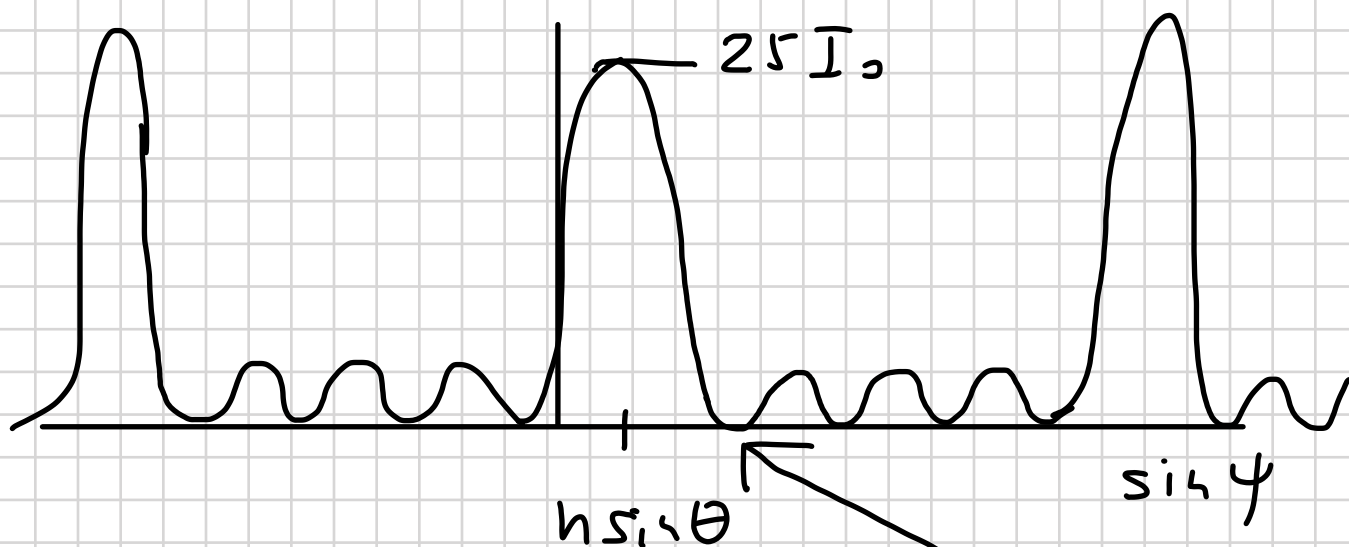
$$\Delta \delta_{\text{before}} = -\frac{2\pi d}{\lambda'} \sin \theta = -\frac{2\pi d n}{\lambda_0} \sin \theta \quad \lambda' = \frac{1}{n} \lambda_0$$

$$\delta_{\text{TOT}} = \frac{2\pi d}{\lambda_0} (\sin \psi - n \sin \theta)$$

c) $I = \sqrt{25} I_0 \frac{\sin^2 \left(\frac{5\pi d}{\lambda_0} (\sin \psi - n \sin \theta) \right)}{\sin^2 \left(\frac{\pi d}{\lambda_0} (\sin \psi - n \sin \theta) \right)}$

Where I_0 is the intensity from single slit.

d)

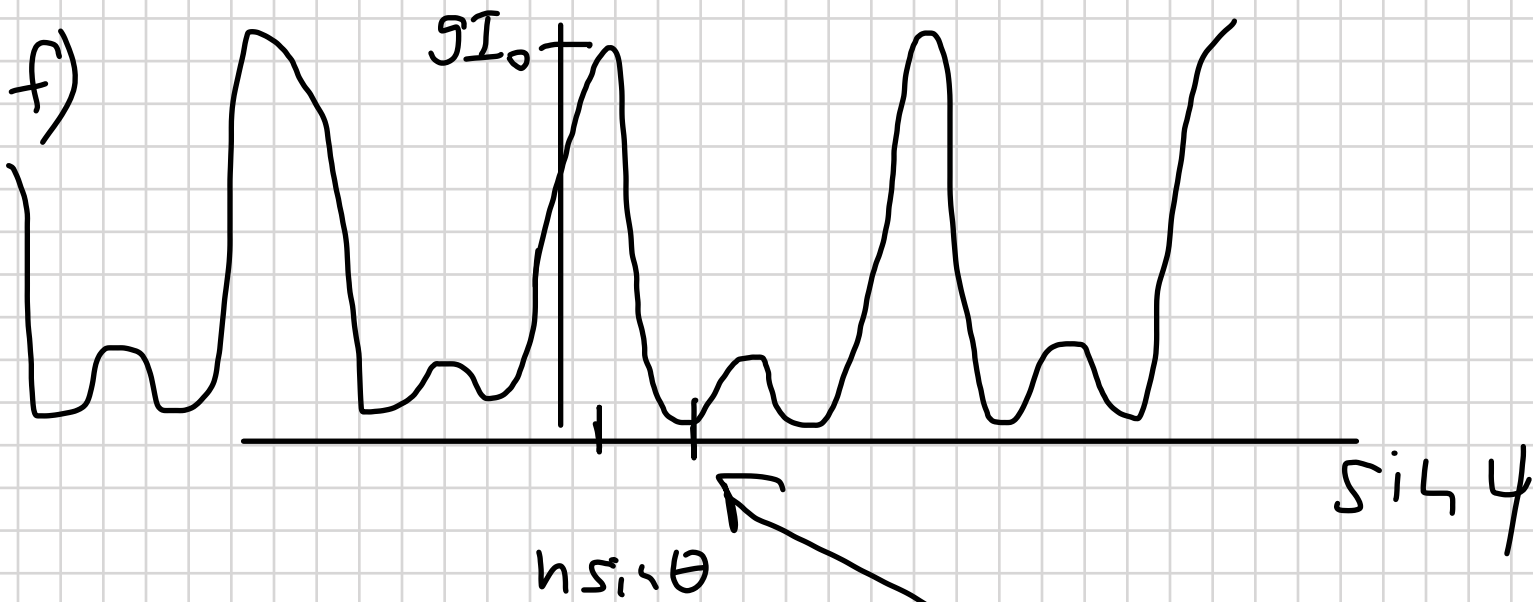


Primary max at $\delta = 0$

1st minimum at $\frac{5\delta}{2} = \pi \quad \sin \psi = n \sin \theta + \frac{\lambda}{5d}$

e) $N \rightarrow 3$

$$I = 9I_0 \frac{\sin^2\left(\frac{3\pi d}{\lambda} (\sin\psi - n\sin\theta)\right)}{\sin^2\left(\frac{\pi d}{\lambda} (\sin\psi - n\sin\theta)\right)}$$



Primary max at $\delta = 0$
 1st minimum at $\frac{3}{2}\delta = \pi$ $\sin\psi = n\sin\theta + \frac{\lambda}{3d}$

Comparison: Max value $25I_0 \leftrightarrow 9I_0$

Primary max pos. Same

1st minimum $+\frac{\lambda}{3d}$ $+\frac{\lambda}{3d}$
 further out for 3
 P.M. wider for 3

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